

# A symmetry problem of elliptic differential operators in potential theory

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## Abstract

This paper is a study of the equation  $(-\Delta_E)^{\frac{\alpha}{2}} u(x) = f(x)$ , where  $(-\Delta_E)^{\frac{\alpha}{2}}$  is an (elliptic pseudo-differential) operator defined by

$$(-\Delta_E)^{-\frac{\alpha}{2}} f = \frac{1}{\Gamma(\frac{\alpha}{2})} \int_0^\infty t^{\frac{\alpha}{2}-1} (H_t * f)(x) dt,$$

$$H_t(x) \equiv H(x, t) = \frac{1}{\sqrt{(4\pi t)^n \eta_1 \eta_2 \cdots \eta_n}} \exp\left(-\sum_i \frac{x_i^2}{4\eta_i t}\right),$$

where  $\eta_1, \eta_2, \dots, \eta_n$  are a set of non-negative numbers that specify the operator. Note that it is an extension of the fractional Laplacian operator  $(-\Delta)^{\frac{\alpha}{2}}$ .

In this paper, we construct a solution, noted as  $J_\alpha f$ , by

$$J_\alpha f(x) \equiv \frac{1}{\beta(\alpha)} \int_{\mathbb{R}^n} \frac{f(y)}{|\eta^{-1} \cdot (x - y)|^{n-\alpha}} dy,$$

where  $|\eta^{-1} \cdot (x - y)|$  is  $\sqrt{\sum_i \eta_i^{-1} (x_i - y_i)^2}$ , and  $\beta(\alpha)^{-1}$  equals

$$\beta(\alpha)^{-1} = \frac{1}{\sqrt{\eta_1 \eta_2 \cdots \eta_n}} \cdot \frac{\Gamma(\frac{n-\alpha}{2})}{\pi^{\frac{n}{2}} 2^\alpha \Gamma(\frac{\alpha}{2})}.$$

Then if we set  $f = \chi_\Omega$  where  $\chi_\Omega$  is the indicator function and  $\Omega$  is some bounded domain in  $\mathbb{R}^n$ , then for all bounded domain  $\Omega$  that is invariant under reflection transformation  $P_m$ , namely  $P_m \Omega = \Omega$  for all  $m = 1, \dots, n$ ,  $J_\alpha f \equiv J_\alpha(x)$  satisfies

$$J_\alpha(x) = J_\alpha(P_m x).$$

The reflection transformation is defined as

$$P_m x = P_m(x_1, \dots, x_m, \dots, x_n) = (x_1, \dots, -x_m, \dots, x_n),$$

where  $m = 1, 2, \dots, n$ .

**摘要:** 在這篇報告中, 我們要探討一個方程式  $(-\Delta_E)^{\frac{\alpha}{2}} u = f$ , 其中  $(-\Delta_E)^{\frac{\alpha}{2}}$  是一個分數次的橢圓形微分算子, 其定義為

$$(-\Delta_E)^{-\frac{\alpha}{2}} f = \frac{1}{\Gamma(\frac{\alpha}{2})} \int_0^\infty t^{\frac{\alpha}{2}-1} (H_t * f)(x) dt,$$

$$H_t(x) \equiv H(x, t) = \frac{1}{\sqrt{(4\pi t)^n \eta_1 \eta_2 \cdots \eta_n}} \exp\left(-\sum_i \frac{x_i^2}{4\eta_i t}\right),$$

其中  $\eta_1, \eta_2, \cdots, \eta_n$  是一群決定其算子特性的參數. 而它是從一般的分數次拉普拉斯算子延伸而得到的.

在報告中, 我們也將找出其一個解, 記為  $J_\alpha f$ , 為

$$J_\alpha f(x) \equiv \frac{1}{\beta(\alpha)} \int_{\mathbb{R}^n} \frac{f(y)}{|\eta \cdot (x - y)|^{n-\alpha}} dy,$$

其中  $|\eta^{-1} \cdot (x - y)|$  代表  $\sqrt{\sum_i \eta_i^{-1} (x_i - y_i)^2}$ , 而  $\beta(\alpha)^{-1}$  等於

$$\beta(\alpha)^{-1} = \frac{1}{\sqrt{\eta_1 \eta_2 \cdots \eta_n}} \cdot \frac{\Gamma(\frac{n-\alpha}{2})}{\pi^{\frac{n}{2}} 2^\alpha \Gamma(\frac{\alpha}{2})}.$$

如果在  $J_\alpha f$  中令  $f = \chi_\Omega$ , 其中  $\chi_\Omega$  是指示函數, 而  $\Omega$  是一個在  $\mathbb{R}^n$  中的有界區域, 則對於所有滿足鏡射變換  $P_m$  的  $\Omega$ , 更精確的說, 對於  $m = 1, \dots, n$ , 都有  $P_m \Omega = \Omega$ ,  $J_\alpha f \equiv J(x)$  滿足

$$J_\alpha(x) = J_\alpha(P_m x).$$

鏡射變換定義為

$$P_m x = P_m(x_1, \cdots, x_m, \cdots, x_n) = (x_1, \cdots, -x_m, \cdots, x_n)$$

其中  $m = 1, 2, \dots, n$ .

## 1 Introduction

The basic idea of this paper is derived from an important concept in potential theory, the Riesz potential  $I_\alpha f$ . It is known that Riesz potential is closely related to the fractional Laplacian operator. It is actually the inverse operator of  $(-\Delta)^{\frac{\alpha}{2}}$ , namely,  $u(x) = I_\alpha f$  if  $(-\Delta)^{\frac{\alpha}{2}} u = f$  [1]. Now we let  $f \equiv \chi_\Omega$ , where  $\chi_\Omega$  is the indicator function. Then this function denoted as  $I_\alpha(x)$  in some bounded domain  $\Omega$  has an interesting property.  $I_\alpha(x)$  is radially symmetric to a center of a ball. In other words,  $u(x)|_{\partial\Omega} = \text{const.}$  if and only if  $\Omega$  is a ball [4].

In this paper, we will extend the fractional Laplacian to an elliptic operator

$$(-\Delta_E)^{\frac{\alpha}{2}} u = \left(-\sum_j \eta_j \frac{\partial^2}{\partial x_j^2}\right)^{\frac{\alpha}{2}} u,$$

where  $\eta_1, \eta_2, \cdots, \eta_n > 0$  and they are independent of the variables. The fractional exponent will be defined in the article. We hope to achieve the following things in the paper:

1. Find the solution of  $(-\Delta_E)^{\frac{\alpha}{2}} u = f$ , which is denoted by  $J_\alpha f(x)$ . Then  $u(x) = J_\alpha f$  if  $(-\Delta_E)^{\frac{\alpha}{2}} u(x) = f(x)$ .

2. Discuss the integrability of  $J_\alpha f$ .
3. Discuss the symmetry property of the solution of  $(-\Delta_E)^{\frac{\alpha}{2}} u = \chi_\Omega$  where  $\Omega$  is an  $n$ -dimensional ellipsoid centered at origin point and axis parallel to the axis  $(x_1, x_2, \dots, x_n)$  of some cartesian coordinate system.
4. Consider symmetry property of the solution of another equation  $(-\Delta_E)^{\frac{\alpha}{2}} u = \chi_\Omega x_i$ , where  $i = 1, 2, \dots, n$ . (The antisymmetric property)

But before doing all this, we will first define some concepts.

## 1.1 Fractional Laplacian

Now we turn to an important concept of this paper: the fractional Laplacian operator  $(-\Delta)^{-\frac{\alpha}{2}}$ . Only the fractional exponent of a positive definite operator can be defined, so we need to take a minus sign in front of the ordinary Laplacian  $\Delta$ .

One way to define  $(-\Delta)^{-\frac{\alpha}{2}}$  is to use the Gamma function  $\Gamma(\alpha)$ . We can start from the fact that for any number  $A$  [1, 3]:

$$A^{-s} = \frac{1}{\Gamma(s)} \int_0^\infty t^{s-1} e^{-tA} dt. \quad (1)$$

If we exchange  $A$  to a Laplacian,  $A \mapsto -\Delta, s \rightarrow \frac{\alpha}{2}$ , then we get the definition.

**Definition 1.** The fractional Laplacian  $(-\Delta)^{-\frac{\alpha}{2}}$  is defined by

$$(-\Delta)^{-\frac{\alpha}{2}} f = \frac{1}{\Gamma(\frac{\alpha}{2})} \int_0^\infty t^{\frac{\alpha}{2}-1} e^{t\Delta} f dt, \quad (2)$$

where

$$e^{\Delta t} f(x) = G_t * f(x) = \int_{\mathbb{R}^n} G_t(x-y) f(y) dy \quad (3)$$

and

$$G(x, t) = G_t(x) = (4\pi t)^{-n} \exp\left(-\frac{|x|^2}{4t}\right) \geq 0. \quad (4)$$

$G_t(x)$  is called the Gauss-Weierstrass kernel [1]. It is the fundamental solution of heat equation, and it is not difficult to see why we use it to define  $e^{t\Delta}$

$$\frac{\partial G_t(x)}{\partial t} = \Delta G_t(x) \iff G_t(x) = e^{\Delta t}, \quad t > 0. \quad (5)$$

However, there is a problem in this definition. When  $\alpha = -2n$ , where  $n$  is a positive integer, then the  $\frac{1}{\Gamma(\frac{\alpha}{2})} = \frac{1}{\Gamma(-n)}$  part will be zero, and the integral part diverges. We fix this problem by taking the limit

$$\lim_{\alpha \rightarrow -2n} \frac{1}{\Gamma(\frac{\alpha}{2})} \int_0^\infty t^{\frac{\alpha}{2}-1} e^{tA} f dt \quad (6)$$

where  $A$  could be any number, and we find this limit to be  $A^n$  by using the equation

$$\frac{\Gamma(s+1)}{A^{-s+1}} = \int_0^\infty t^s e^{-At} dt.$$

So it is reasonable to redefine the fractional Laplacian by taking limits in the definition of it. Now we can define the fractional Laplacian with a positive integer exponent by

$$(-\Delta)^n = \lim_{\alpha \rightarrow 2n} (-\Delta)^{-\frac{\alpha}{2}}. \quad (7)$$

## 1.2 Riesz potential

Riesz potential is closely related to the fractional Laplacian, for it can be seen as an inverse of the fractional Laplacian [1].

**Definition 2.** For any  $n \geq 2$ ,  $0 < \alpha < n$ , and  $x \in \mathbb{R}^n$  the Riesz potential is

$$I_\alpha f(x) = (K_\alpha * f)(x) = \frac{1}{\gamma(\alpha)} \int_{\mathbb{R}^n} \frac{f(y)}{|x-y|^{n-\alpha}} dy, \quad (8)$$

where  $\gamma(\alpha)$  is

$$\gamma(\alpha) = \frac{\pi^{\frac{n}{2}} 2^\alpha \Gamma(\frac{\alpha}{2})}{\Gamma(\frac{n-\alpha}{2})}$$

and

$$K_\alpha = \frac{1}{\gamma(\alpha)} |x|^{\alpha-n} \quad (9)$$

is called the Riesz kernel.

We are going to focus on Riesz potential in a compact domain  $\Omega$  or

$$\frac{1}{\gamma(\alpha)} \int_\Omega \frac{f(y)}{|x-y|^{n-\alpha}} dy = \frac{1}{\gamma(\alpha)} \int_{\mathbb{R}^n} \frac{f(y)}{|x-y|^{n-\alpha}} \chi_\Omega dy, \quad (10)$$

where  $\chi_\Omega$  is the indicator function. The Riesz potential is a singular integral operator, so the concept of integrability is important. In other words, the question will be for  $f \in L^p(\Omega)$ , and  $I_\alpha f \in L^q(\Omega)$ , that  $p, q$  satisfy some condition which makes  $I_\alpha : L^p(\Omega) \rightarrow L^q(\Omega)$  a bounded operator.

This property can be seen by the Hardy-Littlewood-Sobolev inequality [2]:

**Theorem 1.** For  $0 < \alpha < n$ ,  $1 \leq p, q \leq \infty$ ,  $I_\alpha : L^p(\Omega) \rightarrow L^q(\Omega)$  is a bounded operator:

$$\|I_\alpha f\|_q \leq C \|f\|_p, \quad \text{if } \frac{n}{p} \leq \frac{n}{q} + \alpha. \quad (11)$$

*Proof.* See [2]. This theorem says that if  $f \in L^p(\Omega)$ , then for  $x \in \Omega$ ,  $I_\alpha f(x)$  converges absolutely.  $\square$

We are going to see the relationship between fractional Laplacian and the Riesz potential.

**Theorem 2.** For any  $0 < \alpha < n$ , if  $u(x)$  satisfies the equation

$$(-\Delta)^{\frac{\alpha}{2}} u(x) = \rho(x), \quad x \in \mathbb{R}^n, \quad (12)$$

then  $u(x)$  can be written as the convolution of  $K_\alpha$  and  $f$ :

$$u(x) = I_\alpha \rho(x) = (K_\alpha * \rho)(x). \quad (13)$$

*Proof.* The proof is standard [1]. For convenience, we will recall it in the appendix.  $\square$

## 2 Derive $J_\alpha$

### 2.1 Extending the fractional Laplacian

Before extending the fractional Laplacian, we will start by looking at the normal Laplacian first:

$$\Delta \equiv \sum_i \frac{\partial^2}{\partial x_i^2}. \quad (14)$$

We will extend this to

$$-\Delta_E = -\sum_i \eta_i \frac{\partial^2}{\partial \xi_i^2}, \quad \text{where } (\eta_1, \eta_2, \dots, \eta_n > 0), \quad (15)$$

because it is positive definite,  $\eta_1, \eta_2, \eta_3 \dots > 0$ . For the specified case  $\eta_1 = \eta_2 = \dots = \eta_n = 1$ , it reduces to the ordinary Laplacian.

The question is how to define this operator with a fractional exponent  $(-\Delta_E)^{\frac{\alpha}{2}}$ . We can do the same as the original fractional Laplacian:

**Definition 3.** The fractional exponent for the elliptical operator can be written as

$$(-\Delta_E f)^{-\frac{\alpha}{2}} = \frac{1}{\Gamma(\frac{\alpha}{2})} \int_0^\infty t^{\frac{\alpha}{2}-1} e^{\Delta_E t} f dt, \quad (16)$$

where  $e^{\Delta_E t} f = (H * f)(\xi)$ ,  $e^{\Delta_E t} \delta(\xi) \equiv H(t, \xi)$  is the fundamental solution for  $\partial_t u = \Delta_E u$ , and

$$H(\xi, t) = \frac{1}{\sqrt{(4\pi t)^n \eta_1 \eta_2 \dots \eta_n}} \exp\left(-\sum_i \frac{\xi_i^2}{4\eta_i t}\right) \geq 0. \quad (17)$$

(17) can be easily calculated,

$$\frac{\partial H(\xi, t)}{\partial t} - \Delta_E H(\xi, t) = 0 \quad (t > 0, \lim_{t \rightarrow 0} = \delta(x)) \quad (18)$$

and we apply (18) to the Fourier transformation

$$\frac{\partial \hat{H}(\xi, t)}{\partial t} + \left( \sum_i \eta_i \xi_i^2 \right) \hat{H}(\xi, t) = 0, \quad (19)$$

$$\hat{H}(\xi, t) = \exp \left( - \sum_i \eta_i \xi_i^2 t \right) t = 0, \quad (20)$$

$$H(\xi, t) = \prod_i \frac{1}{2\sqrt{\pi t \eta_i}} \exp \left( - \frac{\xi_i^2}{4t \eta_i} \right) = \frac{1}{\sqrt{(4\pi t)^n \eta_1 \eta_2 \cdots \eta_n}} \exp \left( - \sum_i \frac{\xi_i^2}{4\eta_i t} \right). \quad (21)$$

## 2.2 The solution for $\Delta_E$

With all these definitions, we can start to derive the solution for fractional elliptic operator associated to  $\Delta_E$ .

**Theorem 3.** The solution for fractional elliptic operator

$$(-\Delta_E)^{\frac{\alpha}{2}} u(x) = \rho(x) \quad (22)$$

can be taken as  $u(x) = J_\alpha \rho(x)$ , where

$$J_\alpha u(x) \equiv \frac{1}{\beta(\alpha)} \int_{\mathbb{R}^n} \frac{f(y)}{|\eta^{-1} \cdot (x - y)|^{n-\alpha}} dy, \quad (23)$$

and  $|\eta^{-1}(x - y)|$  stands for  $\sqrt{\sum_i \eta_i^{-1}(x_i - y_i)^2}$ , and  $\beta(\alpha)^{-1}$  equals

$$\beta(\alpha)^{-1} = \frac{1}{\sqrt{\eta_1 \eta_2 \cdots \eta_n}} \cdot \frac{\Gamma(\frac{n-\alpha}{2})}{\pi^{\frac{n}{2}} 2^\alpha \Gamma(\frac{\alpha}{2})}. \quad (24)$$

*Proof.* This theorem can be proved by some simple transformation of the variables.

For the equation

$$(-\Delta_E)^{\frac{\alpha}{2}} u(x) = \rho(x), \quad (25)$$

consider a transformation:

$$x_i \mapsto \frac{\xi_i}{\sqrt{\eta_i}}. \quad (26)$$

Then (25) is transformed to

$$(-\Delta)^{\frac{\alpha}{2}} \tilde{u}(\xi) = \tilde{\rho}(\xi). \quad (27)$$

This is just the ordinary fractional Laplacian, so its solution is just the Riesz potential:

$$\tilde{u}(\xi) = \frac{1}{\gamma(\alpha)} \int_{\mathbb{R}^n} \frac{\tilde{\rho}(\zeta)}{|\xi - \zeta|^{n-\alpha}} d\zeta, \quad (28)$$

and we can transform it back to  $x_i$  variable, so the solution will be

$$u(x) = \frac{1}{\sqrt{\eta_1 \eta_2 \cdots \eta_n}} \cdot \frac{\Gamma(\frac{n-\alpha}{2})}{\pi^{\frac{n}{2}} 2^\alpha \Gamma(\frac{\alpha}{2})} \int_{\mathbb{R}^n} \frac{f(y)}{|\eta^{-1} \cdot (x - y)|^{n-\alpha}} dy. \quad (29)$$

□

It is also easy to define the solution for some compact domain  $\Omega$ ; simply set

$$J_\alpha f \equiv \frac{1}{\beta(\alpha)} \int_\Omega \frac{f(y)}{|\eta^{-1}(x-y)|^{n-\alpha}} dy = \frac{1}{\beta(\alpha)} \int_{\mathbb{R}^n} \frac{f(y)\chi_\Omega}{|\eta^{-1}(x-y)|^{n-\alpha}} dy \quad (30)$$

where  $\chi_\Omega$  is the indicator function.

### 3 Integrability

We have to discuss the integrability of  $J_\alpha f$ . Because  $J_\alpha f$  can be turned to  $I_\alpha f$  by changing variables, they should satisfy the same inequality. This has been proven to be true, so we can apply everything in the same way.

**Theorem 4.** Let  $0 \leq q \leq \infty$ ,  $0 < \alpha < n$ . Then  $J_\alpha : L^p(\Omega) \rightarrow L^q(\Omega)$  is a continuous operator

$$\|J_\alpha f\|_{L^q(\Omega)} \leq C \|f\|_{L^p(\Omega)}, \quad \text{for any } \frac{1}{p} \leq \frac{1}{q} + \frac{\alpha}{n}. \quad (31)$$

*Proof.* Before proofing this theorem we need some lemmas.

**Lemma 1.** If a function  $f(x)$  depends only on  $|\eta^{-1}x| \equiv r$  (where the norm stands for  $(\sum_i^n \eta_i^{-1} x_i)^{1/2}$ ), then we have the integral equality

$$\int_{\mathbb{R}^n} f(x) dx = \omega_n \int_0^\infty f(r) r^{n-1} dr, \quad (32)$$

where  $\omega_n$  is

$$\omega_n = \sqrt{\eta_1 \cdots \eta_n} \frac{\pi^{\frac{n}{2}}}{\Gamma(\frac{n}{2} + 1)}. \quad (33)$$

*Proof.* We can start from the fact that [3]

$$\begin{aligned} & \int_{\mathbb{R}_+^n} f(x_1^{b_1} + x_2^{b_2} + \cdots + x_n^{b_n}) x_1^{a_1-1} x_1^{a_2-1} \cdots x_n^{a_n-1} dx \\ &= \frac{\Gamma(\frac{a_1}{b_1}) \Gamma(\frac{a_2}{b_2}) \cdots \Gamma(\frac{a_n}{b_n})}{b_1 \cdots b_n \Gamma(\frac{a_1}{b_1} + \frac{a_2}{b_2} + \cdots + \frac{a_n}{b_n})} \int_0^\infty f(t) t^{\frac{a_1}{b_1} + \frac{a_2}{b_2} + \cdots + \frac{a_n}{b_n} - 1} dt. \end{aligned} \quad (34)$$

$\mathbb{R}_+^n$  is defined as

$$\mathbb{R}_+^n = \{x \in \mathbb{R}^n \mid x_1, \cdots, x_n > 0\}. \quad (35)$$

By setting  $b_1 = b_2 = \cdots = b_n = 2$ , and  $a_1 = a_2 = \cdots = a_n = 1$ , and a transformation,

$$x_i \mapsto \sqrt{\eta_i} x_i, \quad i = 1, 2, \cdots, n, \quad (36)$$

we get

$$\int_{\mathbb{R}_+^n} f(|\eta^{-1}x|^2) dx = \sqrt{\eta_1 \cdots \eta_n} \frac{\pi^{n/2}}{2^n \Gamma(\frac{n}{2})} \int_0^\infty f(t) t^{\frac{n}{2}-1} dt. \quad (37)$$

Last, consider a change of variable  $t = r^2$ :

$$\int_{\mathbb{R}_+^n} f(|\eta^{-1}x|)dx = \sqrt{\eta_1 \cdots \eta_n} \frac{\pi^{n/2}}{2^n \Gamma(\frac{n}{2} + 1)} \int_0^\infty f(r)r^{n-1}dr. \quad (38)$$

By the symmetry of  $f(|\eta^{-1}x|)$ , it is easy to check

$$2^n \int_{\mathbb{R}_+^n} f(|\eta^{-1}x|)dx = \int_{\mathbb{R}^n} f(|\eta^{-1}x|)dx, \quad (39)$$

then the lemma is proven.  $\square$

**Lemma 2.** For some  $1 \leq p, q, r \leq \infty$ , if they satisfy

$$\frac{1}{r} + 1 = \frac{1}{p} + \frac{1}{q}, \quad (40)$$

then

$$\beta(\alpha) \|J_\alpha f\|_r \leq \|f\|_p \|h\|_q, \quad (41)$$

where

$$h(x, y) \equiv h(\eta^{-1}(x - y)) = \frac{1}{|\eta^{-1}(x - y)|^{n-\alpha}}. \quad (42)$$

*Proof.* First, we set

$$\begin{aligned} |J_\alpha f| &= \frac{1}{\beta(\alpha)} \left| \int_{\mathbb{R}^n} f(y)h(x, y)dy \right| \leq \frac{1}{\beta(\alpha)} \int_{\mathbb{R}^n} |h(x, y)f(y)|dy \\ &= \frac{1}{\beta(\alpha)} \int_{\mathbb{R}^n} |f(y)|^{\frac{p}{r}} |f(y)|^{1-\frac{p}{r}} |h(x, y)|^{\frac{q}{r}} |h(x, y)|^{1-\frac{q}{r}} dy. \end{aligned} \quad (43)$$

We can see that,

$$\frac{1}{r} + \left( \frac{1}{p} - \frac{1}{r} \right) + \left( \frac{1}{q} - \frac{1}{r} \right) = \frac{1}{r} + \frac{1}{pr/(p-r)} + \frac{1}{qr/(q-r)} = 1. \quad (44)$$

Then we can apply the Hölder inequality to it:

$$\begin{aligned} \beta(\alpha) |J_\alpha f| &\leq \left( \int_{\mathbb{R}^n} |f(y)|^p |h(x, y)|^q dy \right)^{\frac{1}{r}} \cdot \left( \int_{\mathbb{R}^n} |f(y)|^p dy \right)^{\frac{1}{p} - \frac{1}{r}} \\ &\quad \cdot \left( \int_{\mathbb{R}^n} |h(x, y)|^q dy \right)^{\frac{1}{q} - \frac{1}{r}}. \end{aligned} \quad (45)$$

Take both sides to an exponent  $r$ , then integrate it by  $x$ , and we get

$$\beta(\alpha)^r \|J_\alpha f\|_r^r \leq \left( \int_{\mathbb{R}^n} |f|^p dy \right) \left( \int_{\mathbb{R}^{2n}} |h|^q dx dy \right) \|f\|_p^{r-p} \|h\|_q^{r-q} = \|f\|_p^r \|h\|_q^r \quad (46)$$

and the lemma is proven.  $\square$



**Lemma 3.** For  $n \geq 2, 0 < \alpha < n$ , one has

$$\int_{\Omega} \frac{1}{|\eta^{-1}(x-y)|^{\alpha}} dy \leq \frac{n|E_n|}{n-\alpha} \left( \frac{|\Omega|}{|E_n|} \right)^{1-\frac{\alpha}{n}}, \quad (47)$$

where  $E_n$  is

$$E_n = \sqrt{\eta_1 \cdots \eta_n} \frac{\pi^{n/2}}{\Gamma(\frac{n}{2} + 1)}. \quad (48)$$

It is the volume of a  $n$ -dimensional ellipsoid with axes  $\sqrt{\eta_1}, \sqrt{\eta_2}, \dots, \sqrt{\eta_n}$ .

*Proof.* First, we set  $S \in \mathbb{R}^n$  an  $n$ -dimensional ellipsoid centered at  $x$  with axes  $\sqrt{\eta_1}R, \sqrt{\eta_2}R, \dots, \sqrt{\eta_n}R$  and each parallel to the axis  $x_1, x_2, \dots, x_n$  of coordinate, and  $|S| = E_n R^n$  is the volume of  $S$ . Then we set  $|\Omega| = |S|$ , so that  $R = (|\Omega|/|S|)^{1/n}$ .

$$\int_S \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} = \int_{S \cap \Omega} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} + \int_{S - (S \cap \Omega)} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} \quad (49)$$

and

$$\int_{\Omega} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} = \int_{S \cap \Omega} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} + \int_{\Omega - (S \cap \Omega)} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}}. \quad (50)$$

Because  $S - (S \cap \Omega)$  is inside  $S$ ,  $\frac{1}{|\eta^{-1}(x-y)|^{\alpha}} \leq R^{-\alpha}$ ; therefore

$$\int_{\Omega - (S \cap \Omega)} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} \leq R^{-\alpha} (|\Omega| - |S \cap \Omega|). \quad (51)$$

Similarly,  $\Omega - (S \cap \Omega)$  is outside  $S$ , so that

$$\int_{S - (S \cap \Omega)} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} \geq R^{-\alpha} (|S| - |S \cap \Omega|) = R^{-\alpha} (|\Omega| - |S \cap \Omega|), \quad (52)$$

thus we get

$$\int_{S - (S \cap \Omega)} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} \geq \int_{\Omega - (S \cap \Omega)} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}}, \quad (53)$$

or

$$\begin{aligned} \int_{\Omega} \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} &\leq \int_S \frac{dy}{|\eta^{-1}(x-y)|^{\alpha}} = \int_0^R r^{-\alpha} r^{n-1} |nE_n| dr \\ &= \frac{R^{n-\alpha}}{n-\alpha} nE_n = \frac{n|E_n|}{n-\alpha} \left( \frac{|\Omega|}{|E_n|} \right)^{1-\frac{\alpha}{n}}. \end{aligned} \quad (54)$$

Replace  $\chi_{\Omega} f(x)$  by  $f(x)$  in lemma one in (54), and the lemma is proven.  $\square$

The rest of the proof is obvious. For some  $r$  we can let

$$\frac{1}{r} + \frac{1}{p} = 1 + \frac{1}{q} \quad (55)$$

and if  $1 \leq r \leq 1/(1 - n/\alpha)$  then  $n/p \leq n/q + n$  is satisfied. By Lemma 2

$$\|J_\alpha f\| \leq \beta(\alpha)^{-1} \|h\|_r \|f\|_p. \quad (56)$$

Note that we have replaced  $f(x)\chi_\Omega$  by  $f(x)$ , and  $h(x, y)\chi_\Omega$  by  $h(x, y)$  because (57) only integrates over a bounded domain.

Then by Lemma 3

$$\|h\|_r \leq \left( \frac{nE_n}{n - (n - \alpha)r} \right)^{\frac{1}{r}} \left( \frac{|\Omega|}{|E_n|} \right)^{\frac{1}{r} + \frac{\alpha}{n} - 1}. \quad (57)$$

So the theorem is proven.  $\square$

## 4 The symmetry problem

We know that for  $I_\alpha$  the solution of  $(-\Delta)^{\alpha/2}u = \chi_\Omega$  has some very interesting property, such as its the volume on  $\partial\Omega$  is a constant if any only if  $\Omega$  is a ball [4].

$(-\Delta_E)^{\frac{\alpha}{2}}u = \chi_\Omega$  is invariant under some ‘‘elliptical rotation’’ that preserves  $|\eta^{-1}x|$ , just like the  $(-\Delta)^{\alpha/2}u = \chi_\Omega$  is invariant under rotations that preserve  $|x|$ , but the same property cannot carry over; that is,  $J_\alpha(x)|_\Omega$  will not be a constant where  $\Omega$  is the ellipsoid with axis parallel to the coordinate. It is because not all the transformation that preserves  $|\eta^{-1}x|$  preserves an infinitesimal volume  $dV$  in  $\mathbb{R}^n$ , (if we see this transformation as a coordinate transformation, then it means the Jacobian does not equal one) [3], and therefore  $J_\alpha(x)|_{\partial\Omega}$  does not satisfy this property.

There is a transformation that preserves  $|\eta^{-1}x|$  and infinitesimal volume. It is the reflection transformation  $P_m$  (See Definition 4). It is a discrete transformation, so instead of  $J_\alpha(x)|_{\partial\Omega} = \text{const}$ , we will get  $J_\alpha(x)|_{\partial\Omega} = J_\alpha(P_mx)|_{\partial\Omega}$ . (See Theorem 5.)

**Definition 4.** We are going to introduce the reflection transformation  $P_m : \mathbb{R}^n \rightarrow \mathbb{R}^n$ .

$$P_mx = P_m(x_1, \dots, x_m, \dots, x_n) = (x_1, \dots, -x_m, \dots, x_n), \quad (58)$$

where  $m = 1, 2, \dots, n$ .

For an  $n$ -dimensional ellipsoid with axis  $\sqrt{\eta_1}, \sqrt{\eta_2}, \dots, \sqrt{\eta_n}$ , and each parallel to the axis of the coordinate  $(x_1, x_2, \dots, x_n)$ , which will be noted as  $\Omega$  is symmetric under reflection transformation. That is, for a  $x \in \Omega$ , then  $P_mx \in \Omega$ , and for some  $x \in \partial\Omega$ , then  $P_mx \in \partial\Omega$ .

**Theorem 5.** Let

$$u(x) \equiv \frac{1}{\beta(\alpha)} \int_\Omega \frac{dy}{|\eta^{-1}(x - y)|^{n-\alpha}} = J_\alpha, \quad (59)$$

then there is a property of  $u(x)|_{\partial\Omega}$ .

$$u(x)|_{\partial\Omega} = u(P_mx)|_{\partial\Omega}, \text{ for all } n = 1, 2, \dots, n. \quad (60)$$

For convenience, we need to use a different kind of coordinate instead of the ordinary Cartesian coordinate.

**Definition 5.** We are going to define an elliptical coordinate  $(\rho, \phi_1 \cdots \phi_{n-1})$  with the center at some point  $p$ .

$$\begin{aligned} x_1 - p_1 &= \sqrt{\eta_1} \rho \cos \phi_1 \\ x_2 - p_2 &= \sqrt{\eta_2} \rho \sin \phi_1 \cos \phi_2 \\ &\vdots \\ x_{n-1} - p_{n-1} &= \sqrt{\eta_{n-1}} \rho \sin \phi_1 \cdots \sin \phi_{n-2} \cos \phi_{n-1} \\ x_n - p_n &= \sqrt{\eta_n} \rho \sin \phi_1 \cdots \sin \phi_{n-2} \sin \phi_{n-1} \end{aligned}$$

We can set that  $p \in \partial\Omega$ .

Then another coordinate  $(r, \theta_1, \cdots \theta_{n-1})$  at a point  $p'$ .

$$\begin{aligned} x_1 - p'_1 &= \sqrt{\eta_1} r \cos \theta_1 \\ x_2 - p'_2 &= \sqrt{\eta_2 r \sin \theta_1} \cos \theta_2 \\ &\vdots \\ x_{n-1} - p'_{n-1} &= \sqrt{\eta_{n-1}} r \sin \theta_1 \cdots \sin \theta_{n-2} \cos \theta_{n-1} \\ x_n - p'_n &= \sqrt{\eta_n} r \sin \theta_1 \cdots \sin \theta_{n-2} \sin \theta_{n-1} \end{aligned}$$

Then we can set that  $p' = P_m p \in \partial\Omega$ .

With this coordinate, we shall define a subset in  $\mathbb{R}^n$  by

$$\tau_l^{(k)} = \{x \in \Omega \mid 2^{\frac{l-1}{k}} \leq \rho < 2^{\frac{l}{k}}\}, \quad 1 \leq l \leq k. \quad (61)$$

It is easy to check out that

$$\bigcup_{1 \leq l \leq k} \tau_l^{(k)} = \Omega. \quad (62)$$

We will do the same to coordinate  $(r, \theta_1 \cdots \theta_{n-1})$ .

$$\tau'_l{}^{(k)} = \{x \in \Omega \mid 2^{\frac{l-1}{k}} \leq r < 2^{\frac{l}{k}}\}, \quad 1 \leq l \leq k \quad (63)$$

and

$$\bigcup_{1 \leq l \leq k} \tau'_l{}^{(k)} = \Omega. \quad (64)$$

**Lemma 4.** For any  $k$  and any  $1 \leq l \leq k$ , it satisfies

$$|\tau_l^{(k)}| = |\tau'_l{}^{(k)}|. \quad (65)$$

*Proof.* For any  $x \in \tau_l^{(k)}$  it satisfies the condition

$$\begin{aligned} \left(2\frac{l-1}{k}\right)^2 &\leq \frac{(x_1 - p_1)^2}{\eta_1} + \frac{(x_2 - p_2)^2}{\eta_2} + \dots + \frac{(x_m - p_m)^2}{\eta_m} \\ &+ \dots + \frac{(x_{n-1} - p_{n-1})^2}{\eta_{n-1}} + \frac{(x_n - p_n)^2}{\eta_n} \leq \left(2\frac{l}{k}\right)^2 \end{aligned} \quad (66)$$

and

$$x \in \Omega. \quad (67)$$

If we transform any  $x \in \tau_l^{(k)}$  with the reflection transformation  $P_m$ , then  $P_mx \equiv x'$  satisfies

$$\begin{aligned} \left(2\frac{l-1}{k}\right)^2 &\leq \frac{(P_mx_1 - p_1)^2}{\eta_1} + \frac{(P_mx_2 - p_2)^2}{\eta_2} + \dots + \frac{(P_mx_m - p_m)^2}{\eta_m} \\ &+ \dots + \frac{(P_mx_n - p_n)^2}{\eta_n} \leq \left(2\frac{l}{k}\right)^2 \end{aligned} \quad (68)$$

or

$$\begin{aligned} \left(2\frac{l-1}{k}\right)^2 &\leq \frac{(x_1 - p_1)^2}{\eta_1} + \frac{(x_2 - p_2)^2}{\eta_2} + \dots + \frac{(x_m + p_m)^2}{\eta_m} \\ &+ \dots + \frac{(x_{n-1} - p_{n-1})^2}{\eta_{n-1}} + \frac{(x_n - p_n)^2}{\eta_n} \leq \left(2\frac{l}{k}\right)^2, \end{aligned} \quad (69)$$

and  $x' \in \Omega$ . This is exactly the condition that satisfies for any  $x' \in \tau_l'^{(k)}$ . Thus,

$$P_m\tau_l^{(k)} = \tau_l'^{(k)}. \quad (70)$$

Since the reflection transformation preserves the volume, so that

$$|\tau_l^{(k)}| = |\tau_l'^{(k)}|. \quad (71)$$

□

We know that  $\tau_l^{(k)}$  and  $\tau_l'^{(k)}$  approach to zero as  $k$  approaches to infinity. But how exactly and how rapidly it approaches to zero, we can see it by Lemma 5.

**Lemma 5.** For any integer  $k$ , and some  $1 \leq l \leq k$  the volume of  $\tau_l^{(k)}$  and  $\tau_l'^{(k)}$  satisfy

$$|\tau_l^{(k)}| \leq C(l)k^{-n}, \quad (72)$$

$$|\tau_l'^{(k)}| \leq C(l)k^{-n}, \quad (73)$$

where  $C$  is a constant independent of  $k$  but dependent to  $l$ .

*Proof.* For convenience, we define

$$\sigma_l^{(k)} = \{x \in \mathbb{R}^n \mid 2^{\frac{l-1}{k}} \leq \rho < 2^{\frac{l}{k}}\} \quad (74)$$

and

$$\sigma_l^{\prime(k)} = \{x \in \mathbb{R}^n \mid 2^{\frac{l-1}{k}} \leq r < 2^{\frac{l}{k}}\}. \quad (75)$$

where  $1 \leq l \leq k$ . By the definition of (74) and (75), we can see that  $\tau_l^{(k)} \subseteq \sigma_l^{(k)}$ , and  $\tau_l^{\prime(k)} \subseteq \sigma_l^{\prime(k)}$ ; therefore,  $|\tau_l^{(k)}| \leq |\sigma_l^{(k)}|$ , and  $|\tau_l^{\prime(k)}| \leq |\sigma_l^{\prime(k)}|$ . Since the volume of  $\sigma_l^{(k)}$  and  $\sigma_l^{\prime(k)}$  can be computed

$$|\sigma_l^{(k)}| = \int_{\mathbb{R}^n} \chi_{\sigma_l^{(k)}} dx = \omega_n \int_{2^{\frac{l-1}{k}}}^{2^{\frac{l}{k}}} r^{n-1} dr = \frac{\omega_n}{n-1} \left[ \left(\frac{l}{k}\right)^n - \left(\frac{l-1}{k}\right)^n \right]. \quad (76)$$

We have used Lemma 1 in Equation (76). The  $\omega_n$  has been defined in (33). Therefore, we can see that

$$|\tau_l^{(k)}| \leq |\sigma_l^{(k)}| = \frac{\omega_n}{n-1} \left[ \left(\frac{l}{k}\right)^n - \left(\frac{l-1}{k}\right)^n \right] \equiv C(l)k^{-n}, \quad (77)$$

where  $C(l)$  equals

$$C(l) = \frac{\omega_n}{n-1} [l^n - (l-1)^n]. \quad (78)$$

The case for  $\tau_l^{\prime(k)}$  can be proven in the same way.  $\square$

Now, we can divide the function  $J_\alpha(x)$  into

$$\begin{aligned} J_\alpha(x) &= \frac{1}{\beta(\alpha)} \int_{\Omega} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \\ &= \sum_{1 \leq l \leq k} \int_{\tau_l^{(k)}} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \equiv \sum_{1 \leq l \leq k} j_\alpha^{(l)} \\ &= \sum_{1 \leq l \leq k} \int_{\tau_l^{\prime(k)}} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \equiv \sum_{1 \leq l \leq k} j_\alpha^{\prime(l)}, \end{aligned} \quad (79)$$

where

$$j_\alpha^{(l)}(x) \equiv \int_{\tau_l^{(k)}} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}}, \quad (80)$$

and

$$j_\alpha^{\prime(l)}(x) \equiv \int_{\tau_l^{\prime(k)}} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}}. \quad (81)$$

We are going to define an approximation of  $j^{(l)}$

$$j_{appx.\alpha}^{(l)}(x) = |\tau_l^{(k)}| \frac{1}{|\eta^{-1}(x-y)|^{n-\alpha}}, \quad (82)$$

$$j_{appx.\alpha}^{\prime(l)}(x) = |\tau_l^{\prime(k)}| \frac{1}{|\eta^{-1}(x-y')|^{n-\alpha}}, \quad (83)$$

for some  $y \in \tau_l^{(k)}$  and  $y' \in \tau_l^{\prime(k)}$ .

**Lemma 6.** For any integer  $k$  and  $1 \leq l \leq k$  it satisfies

$$|j_\alpha^{(l)}(p) - j_{\text{approx},\alpha}^{(k)}(p)| \leq C'(l)k^{-\alpha} \quad (84)$$

and

$$|j_\alpha^{\prime(l)}(p') - j_{\text{approx},\alpha}^{\prime(k)}(p')| \leq C'(l)k^{-\alpha}, \quad (85)$$

where  $C'(l)$  is independent of  $k$ .

*Proof.* First, it is obvious that

$$|\tau_l^{(k)}| \min_{x \in \tau_l^{(k)}} \frac{1}{|\eta^{-1}(p-y)|^{n-\alpha}} \leq j_\alpha^{(l)}(p), j_{\text{approx},\alpha}^{(l)}(p) \leq |\tau_l^{(k)}| \max_{x \in \tau_l^{(k)}} \frac{1}{|\eta^{-1}(p-y)|^{n-\alpha}}, \quad (86)$$

where

$$\max_{y \in \tau_l^{(k)}} \frac{1}{|\eta^{-1}(p-y)|^{n-\alpha}} = \frac{1}{(2^{\frac{l-1}{k}})^{n-\alpha}}, \quad (87)$$

and

$$\min_{y \in \tau_l^{(k)}} \frac{1}{|\eta^{-1}(p-y)|^{n-\alpha}} = \frac{1}{(2^{\frac{l}{k}})^{n-\alpha}}. \quad (88)$$

So for a sufficiently large  $k$ , it can satisfy

$$\max_{y \in \tau_l^{(k)}} \frac{|\tau_l^{(k)}|}{|\eta^{-1}(p-y)|^{n-\alpha}} - \min_{x \in \tau_l^{(k)}} \frac{|\tau_l^{(k)}|}{|\eta^{-1}(p-y)|^{n-\alpha}} = \frac{|\tau_l^{(k)}|}{(2^{\frac{l-1}{k}})^{n-\alpha}} - \frac{|\tau_l^{(k)}|}{(2^{\frac{l}{k}})^{n-\alpha}}, \quad (89)$$

so that

$$|j_\alpha^{(l)}(p) - j_{\text{approx},\alpha}^{(k)}(p)| \leq |\tau_l^{(k)}| \left( \frac{1}{(2^{\frac{l-1}{k}})^{n-\alpha}} - \frac{1}{(2^{\frac{l}{k}})^{n-\alpha}} \right) \equiv C'(l)k^{-\alpha}. \quad (90)$$

We have used Lemma 5 in this equation and  $C'(l)$  equals

$$C'(l) = C(l)2^{\alpha-n} \cdot [(l-1)^{\alpha-n} - l^{\alpha-n}]. \quad (91)$$

Then the theorem is proven.  $\square$

Of course, we can basically do the same with  $|\tau_l^{\prime(k)}|$ . Then we can get

$$|j_\alpha^{\prime(l)}(p') - j_{\text{approx},\alpha}^{\prime(k)}(p')| \leq C'(l)k^{-\alpha}. \quad (92)$$

Note that  $C'(l)$  increases as  $l$  increases, and by (91) and (78), we can see that  $C'(l)$  increases in the order of  $k^{n-1} \cdot k^{\alpha-n-1} = k^{\alpha-2}$

Now, back to the main theorem, we can see that  $j_{\text{approx},\alpha}^{(k)}(p) = j_{\text{approx},\alpha}^{\prime(k)}(p')$  because by Lemma 4  $|\tau_l^{(k)}| = |\tau_l^{\prime(k)}|$ , and  $\frac{1}{|\eta^{-1}(p-y)|^{n-\alpha}} = \frac{1}{|\eta^{-1}(p'-y')|^{n-\alpha}}$  for some  $y \in \tau_l^{(k)}$  and  $P_m y = y' \in \tau_l^{\prime(k)}$ , therefore

$$|j_\alpha^{\prime(l)}(p') - j_\alpha^{(l)}(p)| \leq |j_\alpha^{\prime(l)}(p') - j_{\text{approx},\alpha}^{\prime(k)}(p')| + |j_\alpha^{(l)}(p) - j_{\text{approx},\alpha}^{(k)}(p)| \leq 2C'(l)k^{-\alpha}, \quad (93)$$

thus, for sufficiently large  $k$

$$\begin{aligned} |J_\alpha(p) - J_\alpha(p')| &= \left| \sum_l^k j_\alpha^{(l)}(p) - j_\alpha^{(l)}(p') \right| \leq \sum_l^k |j_\alpha^{(l)}(p') - j_\alpha^{(l)}(p)| \\ &\leq \sum_l^k C'(l)k^{-\alpha} \leq kC'(k) \cdot k^{-\alpha}. \end{aligned} \quad (94)$$

Since  $kC'(k)$  increases in the order of  $k^{\alpha-1}$ , so (94) will decrease in the order of  $k^{-1}$  as  $k$  approaches to infinity. So the theorem is proven.  $\square$

## 4.1 Generalization

In Theorem 5, we have assumed  $\Omega$  to be an  $n$ -dimensional ellipsoid centered at the origin point and it has axis of  $\sqrt{\eta_1}, \sqrt{\eta_2}, \dots, \sqrt{\eta_n}$  each parallel to the coordinate  $(x_1 \cdots x_n)$ . But this assumption is superfluous, for all we need is the restriction for  $\Omega$  is  $P_m\Omega = \Omega$  and  $\Omega$  is bounded. From (66) to (70) we can see that Lemma 4 still holds under this restriction, and therefore, so does in Theorem 5.

Another assumption that is superfluous is that we only consider  $p \in \partial\Omega$  and  $P_m p \in \partial\Omega$ . That is, we only consider  $J_\alpha(x)$  under the restriction  $J_\alpha(x)|_{\partial\Omega}$ . We will extend it to any point  $p \in \mathbb{R}^n$  and  $p' = P_m p \in \mathbb{R}^n$ .

We will redefine the coordinate  $(\rho, \phi_1 \cdots \phi_{n-1})$  and  $(r, \theta_1, \cdots, \theta_{n-1})$  in Definition 5 basically in the same way but this time the coordinate will be centered at any point  $p$  and  $P_m p$  which is not necessary on  $\partial\Omega$ .  $|\tau_l^{(k)}|$  and  $|\tau_l'^{(k)}|$  are now written as

$$\tau_l^{(k)} = \{x \in \Omega \mid 2^{\frac{l-1}{k}} \leq \rho < 2^{\frac{l}{k}}\}, k_{min} \leq l \leq k_{max} \quad (95)$$

and

$$\tau_l'^{(k)} = \{x \in \Omega \mid 2^{\frac{l-1}{k}} \leq r < 2^{\frac{l}{k}}\}, k_{min} \leq l \leq k_{max}, \quad (96)$$

where  $k_{max}$  is defined as  $\forall l > k_{max}, \tau_l^{(k)} = \emptyset$ . Since  $P_m \tau_l^{(k)} = \tau_l'^{(k)}$ , so that  $\forall l > k_{max}, \tau_l'^{(k)} = \emptyset$ . Such  $k_{max}$  exists because of the boundedness of  $\Omega$ . Similarly,  $k_{min}$  is defined as  $\forall l < k_{min}, \tau_l^{(k)} = \emptyset$ . If such  $k_{min}$  does not exist, then set  $k_{min} = 1$ .

By this definition, we can get

$$\Omega = \bigcup_{k_{min} \leq l \leq k_{max}} \tau_l^{(k)} = \bigcup_{k_{min} \leq l \leq k_{max}} \tau_l'^{(k)}, \quad (97)$$

and therefore,

$$\begin{aligned}
J_\alpha(x) &= \frac{1}{\beta(\alpha)} \int_\Omega \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \tag{98} \\
&= \sum_{k_{min} \leq l \leq k_{max}} \int_{\tau_l^{(k)}} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \equiv \sum_{k_{min} \leq l \leq k_{max}} j_\alpha^{(l)} \\
&= \sum_{k_{min} \leq l \leq k_{max}} \int_{\tau_l'^{(k)}} \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \equiv \sum_{k_{min} \leq l \leq k_{max}} j_\alpha'^{(l)}.
\end{aligned}$$

The definition of  $j_\alpha^{(l)}(x)$ ,  $j_\alpha'^{(l)}(x)$ ,  $j_{appx,\alpha}^{(l)}(x)$ , and  $j_{appx,\alpha}'^{(l)}(x)$  are still the same, so that we can see Theorem 5 still holds.

The generalization of Theorem 5 is:

**Theorem 6** (generalization). Let

$$u(x) \equiv \frac{1}{\beta(\alpha)} \int_\Omega \frac{dy}{|\eta^{-1}(x-y)|^{n-\alpha}} = J_\alpha. \tag{99}$$

For all bounded domain  $\Omega$  that satisfies  $\Omega = P_m \Omega$

$$u(x) = u(P_m x). \tag{100}$$

## 4.2 The antisymmetric property

In the equation of  $(-\Delta_E)^{\frac{\alpha}{2}} u(x) = f(x)$ , we have set  $f(x) = \chi_\Omega$  and found out that the solution, noted as  $J_\alpha(x)$ , has the symmetric property. Now, if we replace the function  $f(x)$  by another function  $g(x) = \chi_\Omega x_i$ , where  $1 \leq i \leq n$ , then the solution for  $(-\Delta_E)^{\frac{\alpha}{2}} u(x) = g(x)$ , noted as  $J_\alpha g(x) = \mathfrak{J}_\alpha(x)$  will satisfy another property.

**Theorem 7.** Let

$$u(x) \equiv \frac{1}{\beta(\alpha)} \int_\Omega \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}} = \mathfrak{J}_\alpha(x), \tag{101}$$

then there is a property of  $u(x)|_\Omega$ .

$$u(x)|_{\partial\Omega} = u(P_m x)|_{\partial\Omega}, \text{ for } m \neq i, \tag{102}$$

$$u(x)|_{\partial\Omega} = -u(P_m x)|_{\partial\Omega}, \text{ for } m = i. \tag{103}$$

Where  $\Omega$  is an  $n$ -dimensional ellipsoid centered at origin point and axis parallel to the coordinate  $(x_1, x_2 \cdots x_n)$ .

*Proof.* First, we will redefine,  $j_\alpha^{(l)}$ ,  $j_\alpha'^{(l)}$  as

$$j_\alpha^{(l)}(x) \equiv \int_{\tau_l^{(k)}} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \tag{104}$$



and

$$j_{\alpha}^{(l)}(x) \equiv \int_{\tau_l^{(k)}} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}}. \quad (105)$$

Since  $\tau_l^{(k)}$  and  $\tau_l^{\prime(k)}$  are the same as (61) and (63). so that

$$\begin{aligned} \mathfrak{J}_{\alpha}(x) &= \frac{1}{\beta(\alpha)} \int_{\Omega} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \\ &= \sum_{1 \leq l \leq k} \int_{\tau_l^{(k)}} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \equiv \sum_{1 \leq l \leq k} j_{\alpha}^{(l)} \\ &= \sum_{1 \leq l \leq k} \int_{\tau_l^{\prime(k)}} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \equiv \sum_{1 \leq l \leq k} j_{\alpha}^{\prime(l)}. \end{aligned} \quad (106)$$

Now we are going to do something different. Define a subset in  $\tau_l^{(k)}$ ,

$$\pi_{ml}^{(k)} = \{x \in \tau_l^{(k)} \mid \sqrt{\eta_i} \frac{m-1}{k} \leq x_i < \sqrt{\eta_i} \frac{m}{k}\}. \quad (107)$$

Where  $m$  is an integer ranges from  $-k+1$  to  $k$ . Similarly, we define

$$\pi_{ml}^{\prime(k)} = \{x \in \tau_l^{\prime(k)} \mid \sqrt{\eta_i} \frac{m-1}{k} \leq x_i < \sqrt{\eta_i} \frac{m}{k}\}. \quad (108)$$

By (107) and (108), we can see that  $P_m \pi_{ml}^{(k)} = \pi_{ml}^{\prime(k)}$  for  $m \neq i$  and  $P_m \pi_{ml}^{(k)} = \pi_{-m+1l}^{\prime(k)}$  for  $m = i$ , therefore,  $|\pi_{ml}^{(k)}| = |\pi_{ml}^{\prime(k)}|$  for  $m \neq i$ , and  $|\pi_{ml}^{(k)}| = |\pi_{-m+1l}^{\prime(k)}|$  for  $m = i$ .  $\square$

We can see that  $|\pi_{ml}^{(k)}|$  decay to zero as  $k$  approach to infinity, and Lemma 7 told that who rapidly does it approaches to zero.

**Lemma 7.** For any  $-k+1 \leq m \leq k$ ,  $|\pi_{ml}^{(k)}|$  and  $|\pi_{ml}^{\prime(k)}|$  satisfy

$$|\pi_{ml}^{(k)}| \leq Ck^{-n-1} \quad (109)$$

and

$$|\pi_{ml}^{\prime(k)}| \leq Ck^{-n-1}, \quad (110)$$

where  $C$  is a constant independent of  $k$ .

*Proof.* Set

$$|\pi|_{max} \equiv \max\{|\pi_{-k+1l}^{(k)}|, |\pi_{-k+2l}^{(k)}| \cdots |\pi_{k-1l}^{(k)}|, |\pi_{kl}^{(k)}|\} \quad (111)$$

and

$$|\pi|_{min} \equiv \min\{|\pi_{-k+1l}^{(k)}|, |\pi_{-k+2l}^{(k)}| \cdots |\pi_{k-1l}^{(k)}|, |\pi_{kl}^{(k)}|\}. \quad (112)$$

Notice that,

$$2k|\pi|_{min} \leq \sum_{m=-k+1}^k |\pi_{ml}^{(k)}| = |\tau_l^{(k)}|. \quad (113)$$

By Lemma 5 we can see that

$$|\pi|_{\min} \leq \frac{1}{2}C(l)k^{-n-1} \quad (114)$$

or

$$|\pi_{ml}^{(k)}| \leq |\pi|_{\max} \leq \frac{1}{2} \frac{|\pi|_{\max}}{|\pi|_{\min}} C(l)k^{-n-1}. \quad (115)$$

The same can be done with  $|\pi_{ml}^{\prime(k)}|$ , and the theorem is proven.  $\square$

**Lemma 8.** For any integer  $k$  and  $1 \leq l \leq k$  it satisfies

$$|j_{\alpha}^{(l)}(p) + j_{\alpha}^{\prime(l)}(p')| \leq C'(l)k^{-\alpha}. \quad (116)$$

If  $p \in \partial\Omega$ ,  $p' = P_i p \in \partial\Omega$ , and

$$|j_{\alpha}^{(l)}(p) - j_{\alpha}^{\prime(l)}(p')| \leq C'(l)k^{-\alpha}. \quad (117)$$

If  $p \in \partial\Omega$ ,  $p' = P_m p \in \partial\Omega$  where  $m \neq i$ .

The definition of  $j_{\alpha}^{(l)}(x)$  and  $j_{\alpha}^{\prime(l)}(x)$  are in (104) and (105), and  $C'(l)$  is independent of  $k$ .

*Proof.* First, we define

$$j_{\alpha}^{ml}(x) = \int_{\pi_{ml}^{(k)}} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}} \quad (118)$$

and

$$j_{\alpha}^{\prime ml}(x) = \int_{\pi_{ml}^{\prime(k)}} \frac{x_i dy}{|\eta^{-1}(x-y)|^{n-\alpha}}. \quad (119)$$

It is easy to see that

$$j_{\alpha}^{(l)}(x) = \sum_{m=-k+1}^k j_{\alpha}^{ml}(x), \quad \text{and} \quad j_{\alpha}^{\prime(l)}(x) = \sum_{m=-k+1}^k j_{\alpha}^{\prime ml}(x). \quad (120)$$

Consider an approximation of  $j_{\alpha}^{ml}(x)$ , and  $j_{\alpha}^{\prime ml}(x)$

$$j_{\text{appx.}\alpha}^{ml}(x) = |\pi_{ml}^{(k)}| \frac{x_i}{|\eta^{-1}(x-y)|^{n-\alpha}} \quad (121)$$

and

$$j_{\text{appx.}\alpha}^{\prime ml}(x) = |\pi_{ml}^{\prime(k)}| \frac{x_i}{|\eta^{-1}(x-y')|^{n-\alpha}}. \quad (122)$$

For some  $y \in \pi_{ml}^{(k)}$ , and  $y' \in \pi_{ml}^{\prime(k)}$ ,

$$|\pi_{ml}^{(k)}| \min_{x \in \pi_{ml}^{(k)}} \frac{x_i}{|\eta^{-1}(p-y)|^{n-\alpha}} \leq j_{\alpha}^{(ml)}(p), j_{\text{appx.}\alpha}^{(ml)}(p) \leq |\pi_{ml}^{(k)}| \max_{x \in \pi_{ml}^{(k)}} \frac{x_i}{|\eta^{-1}(p-y)|^{n-\alpha}}. \quad (123)$$

By the difinition of  $\pi_{ml}^{(k)}$  and  $\pi_{ml}'^{(k)}$  we can see that

$$\max_{y \in \pi_{ml}^{(k)}} \frac{x_i}{|\eta^{-1}(p-y)|^{n-\alpha}} \leq \frac{m}{k} \frac{1}{(2^{\frac{l-1}{k}})^{n-\alpha}} \quad (124)$$

and

$$\min_{y \in \pi_{ml}'^{(k)}} \frac{x_i}{|\eta^{-1}(p-y)|^{n-\alpha}} \leq \frac{m-1}{k} \frac{1}{(2^{\frac{l}{k}})^{n-\alpha}}. \quad (125)$$

Apply (124) and (125) and Lemma 7 to (123) we can get

$$|j_{\alpha}^{(ml)}(p) - j_{\text{appx.}\alpha}^{(ml)}(p)| \leq C(m, l)k^{-\alpha-1}. \quad (126)$$

Where  $C(m, l)$  equals

$$C(m, l) = 2^{\alpha-n-1} \frac{|\pi|_{\max}}{|\pi|_{\min}} C(l) \left[ \frac{m}{(l-1)^{n-\alpha}} - \frac{m-1}{l^{n-\alpha}} \right]. \quad (127)$$

The can similarly apply the same way to  $j_{\alpha}'^{(ml)}(p')$ , and get

$$|j_{\alpha}'^{(ml)}(p') - j_{\text{appx.}\alpha}'^{(ml)}(p')| \leq C(m, l)k^{-\alpha-1}. \quad (128)$$

On the other hand, by (121) and (122), we can see that  $j_{\text{appx.}\alpha}^{ml}(p) = j_{\text{appx.}\alpha}'^{ml}(p')$  for  $P_m p = p'$  where  $m \neq i$ , and  $j_{\text{appx.}\alpha}^{ml}(p) = -j_{\text{appx.}\alpha}'^{m+1}(p')$  for  $P_i p = p'$ . Therefore, we get

$$\sum_{m=-k+1}^k j_{\text{appx.}\alpha}^{ml}(p) = \sum_{m=-k+1}^k j_{\text{appx.}\alpha}'^{ml}(p') \quad (129)$$

for  $P_m p = p'$ ,  $m \neq i$ .

$$\sum_{m=-k+1}^k j_{\text{appx.}\alpha}^{ml}(p) = - \sum_{m=-k+1}^k j_{\text{appx.}\alpha}'^{ml}(p') \quad (130)$$

for  $P_i p = p'$ .

So for the  $m \neq i$  case, by (120) we can get,

$$\begin{aligned} |j_{\alpha}^{(l)}(p) - j_{\alpha}'^{(l)}(p')| &= \left| \sum_{m=-k+1}^k (j_{\text{appx.}\alpha}^{ml}(p) - j_{\alpha}^{ml}(p)) + \sum_{m=-k+1}^k (j_{\alpha}'^{ml}(p')) - j_{\text{appx.}\alpha}'^{ml}(p') \right| \\ &\leq \sum_{m=-k+1}^k |j_{\text{appx.}\alpha}^{ml}(p) - j_{\alpha}^{ml}(p)| + \sum_{m=-k+1}^k |j_{\text{appx.}\alpha}'^{ml}(p') - j_{\alpha}'^{ml}(p')| \\ &\leq 2 \sum_{m=-k+1}^k C(l, m)k^{\alpha-1} = 2k \cdot C(l)k^{-\alpha-1} \equiv C(l)k^{-\alpha} \end{aligned} \quad (131)$$

and for the  $m = i$  case

$$\begin{aligned}
|j_\alpha^{(l)}(p) + j_\alpha^{(l)}(p')| &= \left| \sum_{m=-k+1}^k (j_\alpha^{ml}(p) - j_{\text{appx},\alpha}^{ml}(p)) + \sum_{m=-k+1}^k (j_\alpha^{ml}(p') - j_{\text{appx},\alpha}^{ml}(p')) \right| \\
&\leq \sum_{m=-k+1}^k |j_{\text{appx},\alpha}^{ml}(p) - j_\alpha^{ml}(p)| + \sum_{m=-k+1}^k |j_{\text{appx},\alpha}^{ml}(p') - j_\alpha^{ml}(p')| \\
&\leq 2 \sum_{m=-k+1}^k C(l, m) k^{\alpha-1} = 2k \cdot C(l) k^{-\alpha-1} \equiv C(l) k^{-\alpha}.
\end{aligned} \tag{132}$$

Note that by (127),  $\sum_{m=-k+1}^k C(l, m)k$  is independent of  $m$ . So we can see that the theorem is proven.  $\square$

We can see that the first equation in Lemma 7 is identical to (93), so by (94) and (106) we can get

$$\mathfrak{J}_\alpha(p)|_{\partial\Omega} = \mathfrak{J}_\alpha(P_m p)|_{\partial\Omega}, \text{ for } m \neq i. \tag{133}$$

By the second equation of Lemma 7, we can see that it is basically the same as (93) but  $j_\alpha^{(l)}(p')$  has been replaced by  $-j_\alpha^{(l)}(p')$ , so by (94) we can see that  $|\mathfrak{J}_\alpha(p) + \mathfrak{J}_\alpha(P_m p)|$  will approach to zero as  $k$  approaches to infinity, so that

$$\mathfrak{J}_\alpha(p)|_{\partial\Omega} = -\mathfrak{J}_\alpha(P_m p)|_{\partial\Omega}, \text{ for } m = i, \tag{134}$$

and therefore, the theorem has been proven.  $\square$

## 5 Appendix [1]

**Theorem 8.** For any  $0 < \alpha < n$ , if  $u(x)$  satisfies the equation

$$-\Delta^{\frac{\alpha}{2}} u(x) = \rho(x), \text{ for any } x \in \mathbb{R}^n \tag{135}$$

then  $u(x)$  can be written as the convolution between  $K_\alpha$  and  $f$

$$u(x) = I_\alpha \rho(x) = (K_\alpha * \rho)(x). \tag{136}$$

*Proof.* We can start by Fourier transformation.

$$\begin{aligned}
\mathcal{F}(-\Delta^{\frac{\alpha}{2}} u(x)) &= -\widehat{\Delta^{\alpha/2} u(\xi)} \\
&= \int_{\mathbb{R}^n} \left( \frac{1}{\Gamma(\frac{-\alpha}{2})} \int_0^\infty t^{-\frac{\alpha}{2}-1} e^{\Delta t} u(x) dt \right) e^{-ix \cdot \xi} dx \\
&= \frac{1}{\Gamma(\frac{-\alpha}{2})} \int_0^\infty t^{-\frac{\alpha}{2}-1} \left( \int_0^\infty G_t * u(x) e^{-ix \cdot \xi} dx \right) dt \\
&= \frac{1}{\Gamma(\frac{-\alpha}{2})} \int_0^\infty t^{-\frac{\alpha}{2}-1} \widehat{G}_t \cdot \widehat{u}(\xi) dt \\
&= \widehat{\rho}(\xi).
\end{aligned} \tag{137}$$

The Fourier transformation of the Gauss kernel was known as

$$\widehat{G}_t(\xi) = (4\pi t)^{-\frac{n}{2}} \int_0^\infty e^{-\frac{|x|^2}{4t}} \cdot e^{-ix \cdot \xi} d\mathbf{x} = e^{-t|\xi|^2}. \quad (138)$$

So,

$$-\widehat{\Delta^{\alpha/2}u}(\xi) = \frac{\widehat{u}(\xi)}{\Gamma(\frac{-\alpha}{2})} \int_0^\infty t^{-\frac{\alpha}{2}-1} e^{-t|\xi|^2} dt = |\xi|^\alpha \widehat{u}(\xi). \quad (139)$$

In this case, we have use the fact that

$$\int_0^\infty t^s e^{-tu^2} dt = \frac{\Gamma(s+1)}{u^{2(s+1)}}. \quad (140)$$

Therefore, we get

$$\widehat{u}(\xi) = |\xi|^{-\alpha} \widehat{\rho}(\xi), \quad u(x) = (2\pi)^{-n} \int_{\mathbb{R}^n} |\xi|^{-\alpha} e^{i\xi \cdot x} d\xi. \quad (141)$$

The solution can be constructed to the convolution between two functions.  $u(x) = (2\pi)^{-n} \phi * \rho(x)$ . Where the  $\phi$  is

$$\phi(x) = \int_{\mathbb{R}^n} |\xi|^{-\alpha} e^{i\xi \cdot x} d\xi, \quad (142)$$

$$\phi(\lambda x) = \int_{\mathbb{R}^n} |\xi|^{-\alpha} e^{i\xi \cdot (\lambda x)} d\xi = \int_{\mathbb{R}^n} |\xi|^{-\alpha} e^{i(\lambda\xi) \cdot x} d\xi. \quad (143)$$

Now consider a transformation:

$$\lambda\xi = \zeta, \quad d\xi = \lambda^{-n} d\zeta.$$

Then,

$$\phi(\lambda x) = \int_{\mathbb{R}^n} |\lambda^{-1}\zeta|^{-\alpha} e^{i\zeta \cdot x} \lambda^{-n} d\zeta = \lambda^{\alpha-n} \int_{\mathbb{R}^n} |\zeta|^{-\alpha} e^{i\zeta \cdot x} d\zeta = \lambda^{\alpha-n} \phi(x). \quad (144)$$

It is obvious that  $\phi(x)$  is a homogeneous function, so we can express  $\phi$  in  $\phi(x) = C|x|^{\alpha-n}$ , and now we are going to determine the constant  $C$ .

$$\begin{aligned} \int_{\mathbb{R}^n} \phi(x) \exp\left(-\frac{|x|^2}{2}\right) dx &= C \int_{\mathbb{R}^n} \exp\left(-\frac{|x|^2}{2}\right) |\xi|^{\alpha-1} dx \\ &= C |\mathbb{S}^{n-1}| \int_0^\infty r^{\alpha-1} \exp\left(-\frac{r^2}{2}\right) dr \\ &= C |\mathbb{S}^{n-1}| 2^{\frac{\alpha}{2}-1} \Gamma\left(\frac{\alpha}{2}\right). \end{aligned} \quad (145)$$

On the other hand,

$$\begin{aligned} \int_{\mathbb{R}^n} \phi(x) \exp\left(-\frac{|x|^2}{2}\right) dx &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |\xi|^{-\alpha} \exp\left(-\frac{|x|^2}{2}\right) e^{i\xi \cdot x} dx d\xi \\ &= (\sqrt{2\pi})^n \int_{\mathbb{R}^n} \exp\left(-\frac{|\xi|^2}{2}\right) |x|^{-\alpha} d\xi \\ &= (\sqrt{2\pi})^n |\mathbb{S}^{n-1}| \int_0^\infty \exp\left(-\frac{r^2}{2}\right) r^{n-\alpha-1} dr \\ &= (\sqrt{2\pi})^n |\mathbb{S}^{n-1}| 2^{\frac{n-\alpha}{2}-1} \Gamma\left(\frac{n-\alpha}{2}\right). \end{aligned} \quad (146)$$

Compare (145) and (146) with the results, we can see that

$$C = (\sqrt{2\pi})^n 2^{-\alpha + \frac{n}{2}} \frac{\Gamma(\frac{\alpha-n}{2})}{\Gamma(\frac{\alpha}{2})}, \quad (147)$$

$$\phi(x) = (\sqrt{2\pi})^n 2^{-\alpha + \frac{n}{2}} \frac{\Gamma(\frac{\alpha-n}{2})}{\Gamma(\frac{\alpha}{2})} |x|^{\alpha-n}. \quad (148)$$

And finally, we get our solution

$$u(x) = (2\pi)^{-n} \phi * \rho(x) = \pi^{-\frac{n}{2}} 2^{-\alpha} \frac{\Gamma(\frac{\alpha-n}{2})}{\Gamma(\frac{\alpha}{2})} \int_{\mathbb{R}^n} \frac{\rho(y)}{|x-y|^{n-\alpha}} = I_\alpha \rho(x) dx. \quad (149)$$

If the exponent  $\alpha = 2n$  ( $n$  is a natural number), then by definition:

$$-\widehat{\Delta^n u}(\xi) = \lim_{\alpha \rightarrow 2n} -\widehat{\Delta^{\alpha/2} u}(\xi) = \lim_{\alpha \rightarrow 2n} \frac{1}{\Gamma(\frac{-\alpha}{2})} \int_0^\infty t^{-\frac{\alpha}{2}-1} \widehat{G}_t \cdot \widehat{u}(\xi) dt. \quad (150)$$

By calculation, this limit will be

$$\lim_{\alpha \rightarrow 2n} -\widehat{\Delta^{\alpha/2} u}(\xi) = |\xi|^\alpha \widehat{u}(\xi). \quad (151)$$

□

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## References

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