

Global Geometry of Surfaces

(Manuscripts by Chin-Lung Wang)

These pages are selected from the second half of the course notes based on DoCarmo's book, with some supplementary results.

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Gauss equation (Theorema Egregium)

"a most excellent theorem"

$$K = \frac{eg - f^2}{EG - F^2} \quad \text{how to get } e, g, f$$

$$e = N \cdot X_{uu} = \frac{|X_u X_v X_{uu}|}{\sqrt{EG - F^2}}$$

$$f = N \cdot X_{uv} = \frac{|X_u X_v X_{uv}|}{\sqrt{EG - F^2}}$$

$$g = N \cdot X_{vv} = \frac{|X_u X_v X_{vv}|}{\sqrt{EG - F^2}}$$

$$eg = \frac{1}{EG - F^2} \begin{pmatrix} \boxed{X_u X_u} & X_u X_v & X_u X_{uu} \\ X_v X_u & \boxed{X_v X_v} & X_v X_{uv} \\ X_{uu} X_u & X_{uv} X_v & \boxed{X_{vv} X_v} \end{pmatrix}$$

$$= \begin{vmatrix} E & F & F_v - \frac{1}{2} G_u \\ F & G & \frac{1}{2} G_v \\ \frac{1}{2} E_u & F_u - \frac{1}{2} E_v & ? \end{vmatrix}$$

$$\begin{aligned} X_u \cdot X_{vv} &= (X_u \cdot X_v)_v - X_{uv} \cdot X_v \\ &= F_v - \frac{1}{2} (X_v \cdot X_v)_u = F_v - \frac{1}{2} G_u \end{aligned}$$

$$\begin{aligned} X_{uu} \cdot X_v &= (X_u \cdot X_v)_u - X_u \cdot X_{uv} \\ &= F_u - \frac{1}{2} (X_u \cdot X_u)_v = F_u - \frac{1}{2} E_v \end{aligned}$$

$$X_{uu} \cdot X_{vv} = (X_u \cdot X_{vv})_u - X_u \cdot X_{vvu} = \frac{1}{2} E_v$$

$$= (F_v - \frac{1}{2} G_u)_u - (X_u \cdot X_{uv})_v + X_{uv} \cdot X_{uv} \dots ?$$

only get $X_{uu} \cdot X_{vv} - X_{uv} \cdot X_{uv} = (F_v - \frac{1}{2} G_u)_u - \frac{1}{2} E_{vv}$

but $f^2 = \frac{1}{EG - F^2} \begin{vmatrix} \boxed{X_u X_u} & X_u X_v & X_u X_{uv} \\ X_v X_u & \boxed{X_v X_v} & X_v X_{uv} \\ X_{uv} X_u & X_{uv} X_v & \boxed{X_{uv} X_{uv}} \end{vmatrix}$ By combining terms

$$\text{get } K = \frac{1}{(EG - F^2)^2} \left\{ \begin{vmatrix} E & F & F_v - \frac{1}{2} G_u \\ F & G & \frac{1}{2} G_v \\ \frac{1}{2} E_u & F_u - \frac{1}{2} E_v & F_{uv} - \frac{1}{2} G_{uu} - \frac{1}{2} E_{vv} \end{vmatrix} - \begin{vmatrix} E & F & \frac{1}{2} E_v \\ F & G & \frac{1}{2} G_u \\ \frac{1}{2} E_v & \frac{1}{2} G_u & 0 \end{vmatrix} \right\}$$

conclusion: K depends only on ds^2 up to 2nd order.

Isometric Problem of Surface in \mathbb{R}^3 can be transformed into

$$\begin{cases} \alpha_v = p\beta \\ \beta_u = q\alpha \end{cases} \quad \begin{aligned} d_{vu} &= p_u\beta + p\beta_u \\ &= p_u\beta + pq\alpha = (\log p)_u \beta + p q \alpha \end{aligned}$$

or, in graph can be of

"Monge-Ampere Equation" possibly a degenerate hyp. equation!

* Codazzi 方程

compatibility equations

standard frame X_u, X_v, N

$$X_u = X_1$$

$$X_v = X_2$$

Equation (运动方程)

$$(X_u)_u =$$

$$X_{11} = \Gamma_{11}^1 X_1 + \Gamma_{11}^2 X_2 + eN$$

$$X_{12} = X_{21} = \Gamma_{12}^1 X_1 + \Gamma_{12}^2 X_2 + fN$$

$$X_{22} = \Gamma_{22}^1 X_1 + \Gamma_{22}^2 X_2 + gN$$

observation: All Γ_{ij}^k can be determined by E, F, G and their 1st derivatives

$$N_u = \begin{cases} N_1 = a_{11} X_1 + a_{12} X_2 \\ N_2 = a_{21} X_1 + a_{22} X_2 \end{cases}$$

compatibility:

$$X_{112} = X_{121} : \quad \partial_2 \Gamma_{11}^1 X_1 + \Gamma_{11}^1 X_{12} + \partial_2 \Gamma_{11}^2 X_2 + \Gamma_{11}^2 X_{22} + e_2 N + e N_2$$

$$= \partial_1 \Gamma_{12}^1 X_1 + \Gamma_{12}^1 X_{11} + \partial_1 \Gamma_{12}^2 X_2 + \Gamma_{12}^2 X_{21} + f_1 N + f N_1$$

Gauss - Codazzi 方程

① compare X_1 with

$$\begin{aligned} \partial_2 \Gamma_{11}^1 + \Gamma_{11}^1 \Gamma_{12}^1 + \Gamma_{11}^2 \Gamma_{12}^1 + e a_{21} \\ = \partial_1 \Gamma_{12}^1 + \Gamma_{12}^1 \Gamma_{11}^1 + \Gamma_{12}^2 \Gamma_{12}^1 + f a_{11} \end{aligned}$$

had seen: $a_{11} = \frac{fF - eG}{EG - F^2}$; $a_{21} = \frac{gF - fG}{EG - F^2}$

$$\begin{aligned} e a_{21} - f a_{11} &= \frac{1}{EG - F^2} (e g F - f f G - f^2 F + f e G) \\ &= \frac{eg - f^2}{EG - F^2} F = F \cdot K \end{aligned}$$

get $F \cdot K = \partial_1 \Gamma_{12}^1 - \partial_2 \Gamma_{11}^1 + \Gamma_{12}^1 \Gamma_{12}^2 - \Gamma_{11}^2 \Gamma_{22}^1$

This also proves the Theorema Egregium.

Remark: compare X_2 with get same thing. (no e, f, g involved)

② compare N with: (from $(X_{11})_2 = (X_{12})_1$:)

$$\Gamma_{11}^1 f + \Gamma_{11}^2 g + e_2 = \Gamma_{12}^1 e + \Gamma_{12}^2 f + f_1$$

i.e. $e_2 - f_1 = -\Gamma_{12}^1 e + (\Gamma_{12}^2 - \Gamma_{11}^1) f - \Gamma_{11}^2 g$

compare N from $(X_{12})_2 = (X_{22})_1$:

$$f_2 - g_1 = \Gamma_{22}^1 e + (\Gamma_{22}^2 - \Gamma_{12}^1) f - \Gamma_{12}^2 g$$

(this totally symmetric via $1 \leftrightarrow 2, f \leftrightarrow f, e \leftrightarrow g$)

Conclusion:

① Gauss' equation gives intrinsic char of K

② Codazzi equation gives the necessary condition (given I) for a given II being possible for a surface $S \hookrightarrow \mathbb{R}^3$. (compatibility equations)

• Theorem (Bonnet): Local ^{existence &} Uniqueness of surfaces in \mathbb{R}^3 .

Student of Riemann.

\leftrightarrow Gauss - Codazzi Equation

Proof of Bonnet's Theorem:

$$\frac{\partial x^k}{\partial u^\alpha} = U^\alpha_k(u, x)$$

$k = 1 \dots n$; number of eqn's
 $\alpha = 1 \dots m$ = $n \cdot m$

number of functions = " k " m

compatibility conditions:

$$\frac{\partial^2 x^k}{\partial u^\alpha \partial u^\beta} = "U^\alpha_{\beta}{}^k" = U^\alpha_{\beta}{}^k + U^\alpha_{\gamma}{}^k \frac{\partial x^j}{\partial u^\beta} \Big|_{\substack{\text{index} \\ \alpha, \beta \dots}} \Big|_{\substack{\text{index} \\ i, j \dots}} U^\alpha_k(u_1, \dots, u_m, x_1, \dots, x_n)$$

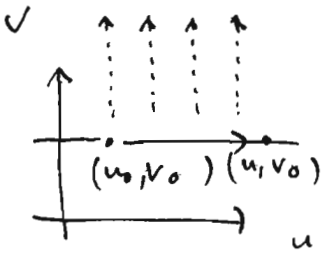
ie. require

$$* \quad U^\alpha_{\beta}{}^k + \sum_j U^\alpha_{\gamma}{}^k \frac{\partial x^j}{\partial u^\beta} = U^\alpha_{\beta}{}^k + \sum_j U^\alpha_{\gamma}{}^k \frac{\partial x^j}{\partial u^\alpha} \quad \text{rep by } U^\alpha_{\beta}{}^k$$

imp. by $U^\alpha_{\beta}{}^k$

pf: let $m=2$. (then use induction)

$$* \quad \left\{ \begin{array}{l} \frac{\partial x^k}{\partial u} = U^k(u, v, x) \\ \frac{\partial x^k}{\partial v} = V^k(u, v, x) \end{array} \right.$$



① fix v_0 , varies u

solve $\frac{dy^k}{du} = U^k(u, v_0, y)$

with initial condition $y^k(u_0, v_0) = x^k_0$

this is just an ODE, always solvable.

② Take y^k as initial data, solve for each u

$$\frac{dx^k}{dv} = V^k(u, v, x) \quad \text{initial data } x^k(u, v_0) = y^k(u, v_0)$$

Need to check that $\frac{\partial x^k}{\partial u} = U^k(u, v, x)$ **

$$\left(\frac{\partial x^k}{\partial u} \right)_{v=v_0} = \frac{dy^k}{du} = U^k(u, v_0, x) \quad \text{is ok.}$$

To check **, need only verify both LHS, RHS satisfies the same ODE. (in variable v)

For this purpose we need the solved $x(u, v)$ to be at least C^2 , say, let U^k be C^1 .

A. $\frac{\partial}{\partial v} \left(\frac{\partial x^k}{\partial u} \right) = \frac{\partial}{\partial u} \left(\frac{\partial x^k}{\partial v} \right) = \frac{\partial}{\partial u} \left(v^k(u, v, x) \right)$
 the direction

B. $\frac{\partial}{\partial v} \left(u^k(u, v, x) \right)$ another direction

A : $\frac{\partial}{\partial v} \left(\frac{\partial x^k}{\partial u} \right) = \frac{\partial v^k}{\partial u} + \sum_i \frac{\partial v^k}{\partial x^i} \left(\frac{\partial x^i}{\partial u} \right)$

B : $\frac{\partial}{\partial v} u^k = \frac{\partial u^k}{\partial v} + \sum_i \frac{\partial u^k}{\partial x^i} \cdot \frac{\partial x^i}{\partial v}$

done!

by construction
 this is already v^i
 by compatibility

Applications to Bonnet's thm:

Frame $\{x_u, x_v, N\}$

get $\frac{\partial v^k}{\partial u} + \sum_i \frac{\partial v^k}{\partial x^i} \cdot u^i$ *

fix an initial point and the compatibility condition turns out to be merely the Codazzi equation

Reason: View it as nine functions $F^i, i=1, \dots, 9$

the PDE is then

$$\begin{aligned} \frac{\partial F^i}{\partial u} &= \\ \frac{\partial F^i}{\partial v} &= \end{aligned} \left(\begin{array}{l} \text{here are all functions} \\ \text{of } x_u, x_v \text{ (not } x) \\ \text{(indeed } E, F, G) \text{ and } e, f, g \end{array} \right)$$

So the compatibility condition is

$$\frac{\partial F^i}{\partial u \partial v} = \frac{\partial F^i}{\partial v \partial u} \quad \#$$

• To solve x , simply take

$$x(u, v) = \int_{u_0}^u x_1(\xi, v_0) d\xi + \int_{v_0}^v x_2(u_0, \eta) d\eta + x(u_0, v_0)$$

Then to check it really satisfies the Codazzi equation.

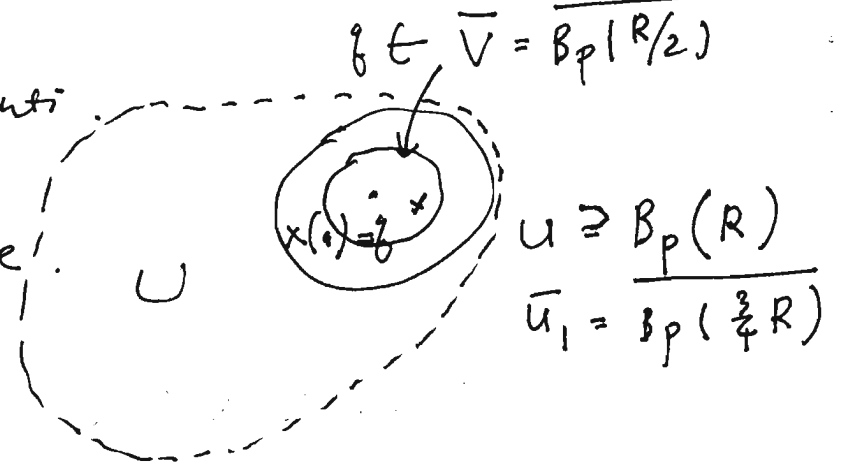
• The uniqueness may be reduced to the uniqueness of curves. eg. for $C = x(u(t), v(t))$ get

$$\begin{cases} x' = x_1 u' + x_2 v' = R(x_1, x_2, t) \\ x_1' = x_{11} u' + x_{12} v' = S(x_1, x_2, t) \\ x_2' = x_{21} u' + x_{22} v' = T(x_1, x_2, t) \end{cases}$$

via Codazzi eqn.

So $(x_1(t), x_2(t))$
 are uniquely
 det. by initial &
 E, F, G, e, f, g *

(Global) Existence and conti of flow $X(\xi, t)$ in the Lipschitz case.



Only need to prove the local case.

Setup: ODE on \mathbb{R}^n

$$X' = F(X) \quad F: U \rightarrow \mathbb{R}^n \text{ Lipschitz}$$

$$|F| \leq M$$

$$\Leftrightarrow x(t) = x(a) + \int_a^t F(x(u)) du$$

$\underbrace{\hspace{10em}}_{\varphi_\xi(x)}$

$$|\varphi_\xi(x) - \varphi_\xi(y)| \leq \left| \int_a^t |F(x(u)) - F(y(u))| du \right|$$

$$\leq \underbrace{|t-a| \cdot M}_{u \in J} \sup_{u \in J} |x(u) - y(u)|$$

$C^0(J, U)$
 $J = [a-d, a+d]$

① make $\delta \cdot M \leq \frac{1}{2} \Rightarrow |\varphi_\xi(x) - \varphi_\xi(y)|_{C^0} \leq \frac{1}{2} |x - y|_{C^0}$

② make $\delta \cdot M \leq \frac{R}{4} \Rightarrow |x(t)| \leq |\xi| + \frac{R}{4} \leq \frac{3}{4} R$

$\Rightarrow \varphi_\xi: C^0(J, \bar{U}_1) \rightarrow C^0(J, \bar{U}_1)$ cont. mapping

\hookrightarrow complete space $C^0(J, \bar{U}_1)$ w.r.t

$\Rightarrow \exists!$ fixed pt X_ξ in fact given by $\varphi_\xi^n(x)$ as $n \rightarrow \infty$ for any x .

Lipschitz conti of X_ξ in ξ :

$$|\varphi_{\xi_1}(X_\xi) - X_\xi| = |\xi_1 - \xi| \quad \forall n$$

$$\Rightarrow |\varphi_{\xi_1}^n(X_\xi) - X_\xi| \leq \sum_{i=1}^n |\varphi_{\xi_1}^i(X_\xi) - \varphi_{\xi_1}^{i-1}(X_\xi)| < 2|\xi_1 - \xi|$$

$n \rightarrow \infty \quad |X_{\xi_1} - X_\xi| \leq 2|\xi_1 - \xi| \quad |X(\xi_1, t) - X(\xi, t)|$ use Δ -ineq. *

ODE: $x' = f(x)$ st. $f \in C^1(0, \mathbb{R}^n)$; $\phi(t, x)$ flow $\exists! C^0$.

let $x(t)$ parti sol. in $J \ni 0$, $x_0 = x(0)$
closed int.

Variational equation:

let $A(t) = DF_x(t) \in C^0$, consider $u' = A(t)u$ (*)

idea: if $u_0 = u(0)$ small, $x(t) + u(t)$ approx. sol with initial data $x_0 + u_0$.

let $u(t, \xi)$ be the flow of VE, i.e. $u(0, \xi) = \xi \quad \forall \xi \in \mathbb{R}^n$ &

$y(t, \xi)$ " of $x' = f(x)$, $y(0, \xi) = x_0 + \xi \quad x_0 + \xi \in \mathcal{O}$

Prop. $\lim_{\xi \rightarrow 0} \frac{|y(t, \xi) - x(t) - u(t, \xi)|}{|\xi|} = 0$ unif on $t \in J$.

Thm: Smoothness of the flow. (Kirsch-Smale): $\phi \in C^1$.

pf: To get $\frac{\partial \phi}{\partial x} \equiv D_2 \phi$, for ξ small

$$\phi(t, x_0 + \xi) - \phi(t, x_0) = y(t, \xi) - x(t) = u(t, \xi) + o(|\xi|)$$

$$\Rightarrow D_2 \phi(t, x_0) \xi = u(t, \xi) \text{ lin in } \xi \text{ w.r.t. } t, x_0 \text{ flows } \square$$

So easy! where is the dependence?

back: Back to (*) get $\frac{d}{dt} D_2 \phi(t, x_0) = DF_x \phi(t, x_0) \cdot D_2 \phi(t, x_0)$

with initial $D_2 \phi(0, x_0) = \text{id}_{\mathbb{R}^n}$. Hence

pf of prop: Integral equations:

$F \in C^k \Rightarrow \phi \in C^k$
by induction.

$$x(t) = x_0 + \int_0^t f(x(s)) ds$$

$$y(t, \xi) = x_0 + \xi + \int_0^t f(y(s, \xi)) ds$$

$$u(t, \xi) = \xi + \int_0^t DF_x(s) (u(s, \xi)) ds$$

$$\Rightarrow |y(t, \xi) - x(t) - u(t, \xi)| \leq \int_0^t |f(y(s, \xi)) - f(x(s)) - DF_x(s) u(s, \xi)| ds$$

$$\stackrel{\text{say}}{t \geq 0} \leq \int_0^t |DF_x(s) (y(s, \xi) - x(s) - u(s, \xi))| ds + \int_0^t |R(x(s); y(s, \xi) - x(s))| ds$$

Taylor

$$\Rightarrow g(t) := \text{LHS} \leq \max_{s \in J} |DF_x(s)| \int_0^t g(s) ds + \dots \quad \epsilon := \text{any } > 0$$

$\epsilon = |y(s, \xi) - x(s)|$

(*) for $|y| \leq \delta_0, s \in J$.

such a $\delta_0 > 0$ exists by

Ex. Gronwall's inequality: if $u(t) \leq C + \int_0^t K u(s) ds$
with $u \in C([0, a])$, $u \geq 0$, then $u(t) \leq C e^{Kt}$

in (*), $C = |\xi|$ can be taken small: $|y(s, \xi) - x(s)| \leq |\xi| e^{Ks} \leq \delta_0$

$$\Rightarrow g(t) \leq N \int_0^t g(s) ds + \epsilon |\xi| \cdot C_1(K, |J|) \Rightarrow g(t) \leq C \cdot \epsilon \cdot |\xi| e^{Nt} \quad **$$

Covariant diff. Parallel transl. & geodesics (intrinsic notion)

$$V = aX_1 + bX_2$$

$$\frac{DV}{dt} := a'X_1 + aX_{11}u' + aX_{12}v' + b'X_2 + bX_{21}u' + bX_{22}v' \quad \text{proj to TS}$$

$$= (a' + \Gamma_{11}^1 u'a + \Gamma_{12}^1 v'a + \Gamma_{21}^1 u'b + \Gamma_{22}^1 v'b) X_1 + (b' + \Gamma_{11}^2 u'a + \Gamma_{12}^2 v'a + \Gamma_{21}^2 u'b + \Gamma_{22}^2 v'b) X_2$$

when $(u,v)(t)$ is given, this is simply linear ODE.

where $V = \alpha' = u'X_1 + v'X_2$ system of
 get not nec. by arc length, a priori

$$\frac{D\alpha'}{dt} = (u'' + \Gamma_{11}^1 u'^2 + 2\Gamma_{12}^1 u'v' + \Gamma_{22}^1 v'^2) X_1 + (v'' + \Gamma_{11}^2 u'^2 + 2\Gamma_{12}^2 u'v' + \Gamma_{22}^2 v'^2) X_2$$

geodesics $\equiv D\alpha'/dt = 0 \Rightarrow |\alpha'| = \text{const} \Rightarrow t \sim \text{arc length}$

examples of parallel transl.

③ surface of revolution:

$$X(u,v) = (f(v) \cos u, f(v) \sin u, g(v))$$

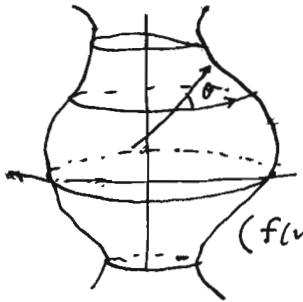


$$\frac{4\pi}{8} = \frac{\pi}{2}!$$

$$A = -\pi + (CA + CB + CC)$$



helix = line



$(f(v), g(v))$

$$\begin{cases} E = f^2, F = 0, G = f'^2 + g'^2 \\ E_1 = 0, E_2 = 2ff'; F_i = 0 \\ G_1 = 0, G_2 = 2(f'f'' + g'g'') \end{cases}$$

$$\begin{aligned} \Gamma_{11}^1 &= 0, \Gamma_{11}^2 = -\frac{ff'}{f'^2 + g'^2} \\ \Gamma_{12}^1 &= \frac{f'f''}{f^2}, \Gamma_{12}^2 = 0 \\ \Gamma_{22}^1 &= 0, \Gamma_{22}^2 = \frac{f'f'' + g'g''}{f'^2 + g'^2} \end{aligned}$$

$$\text{ie. } \begin{cases} u'' + \frac{2ff'}{f^2} u'v' = 0 \\ v'' - \frac{ff'}{f'^2 + g'^2} u'^2 + \frac{f'f'' + g'g''}{f'^2 + g'^2} v'^2 = 0 \end{cases}$$

equiv. under I.

$$\text{ie. } ds^2 = f^2 du^2 + (f'^2 + g'^2) dv^2$$

$$| = f^2 \left(\frac{du}{ds}\right)^2 + (f'^2 + g'^2) \left(\frac{dv}{ds}\right)^2$$

$$\Rightarrow (f^2 u')' = 2ff' u'v' + f^2 u'' = 0$$

ie. $f^2 u' = \text{const } C$ (" is wrt arc length s)
 Clairaut's relation in $[u,v]$

$$\frac{dv}{ds} = \frac{C}{f \sqrt{f^2 - C^2}}$$

$$f' \cos \theta = C; \cos \theta = \frac{\langle X_1, X_1 u' + X_2 v' \rangle}{|X_1|} = f u'$$

$$f^2 \frac{dv}{ds} = C \Rightarrow f^2 \frac{dv}{ds} \sqrt{1 - \frac{C^2}{f^2}} = C$$

get $u = \int \dots dv$

$$\frac{\partial X_i}{\partial u} = (X_{ii})^T = (\Gamma_{ii}^1 X_1 + \Gamma_{ii}^2 X_2 + e_N)^T$$

$$= \Gamma_{ii}^1 X_1 + \Gamma_{ii}^2 X_2$$

$$D_i X_j = X_{ij}^T = \sum_k \Gamma_{ij}^k X_k \quad \text{covariant differentiation}$$

$$g_{kl} = X_k \cdot X_l$$

$$\partial_i g_{kl} = X_{ki} \cdot X_l + X_k \cdot X_{li}$$

$$= \Gamma_{ki}^j g_{jl} + \Gamma_{li}^j g_{jk}$$

$$X_l \cdot X_{ij}^T = \sum_k \Gamma_{ij}^k g_{kl}$$

$$X_l \cdot X_{ij}$$

$$(X_l \cdot X_i)_j - X_{lj} \cdot X_i$$

$$\partial_j g_{li} - \Gamma_{lj}^k g_{ki}$$

General formula.

$$\text{ie. } \partial_j g_{li} = \Gamma_{ij}^k g_{kl} + \Gamma_{lj}^k g_{ki} \xrightarrow{\text{cyclic trick}} \Gamma_{lj}^k g_{ki} \quad \text{formula for } \Gamma_{ij}^k$$

Holonomy point of view:

$$\frac{1}{2} \frac{1}{\sqrt{FG}} \left(-E_2 \frac{du}{dt} + G_1 \frac{dv}{dt} \right)$$

① on curve with $F=0$.

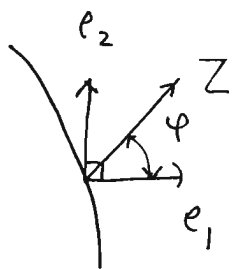
$$\downarrow \left[\frac{D\vec{e}_1}{dt} \right]$$

Algebraic value

$$= \frac{D\vec{e}_1}{dt} \cdot \vec{e}_2$$

what is it really?

"this is indeed a moving frame of"



$$\varphi := \angle(\vec{e}_1, Z)$$

$$\cos \varphi = Z \cdot e_1$$

$$-\varphi' \sin \varphi = Z' \cdot e_1 + Z \cdot e_1'$$

"parallel"

$$Z \cdot \lambda e_2$$

$$\lambda \cos\left(\frac{\pi}{2} - \varphi\right)$$

$$= \lambda \sin \varphi$$

②

$$\text{ie. } \lambda = -\varphi'$$

True for arbitrary unit v.f e_1 .

cor: This \Rightarrow for any 2 unit v.f w, v

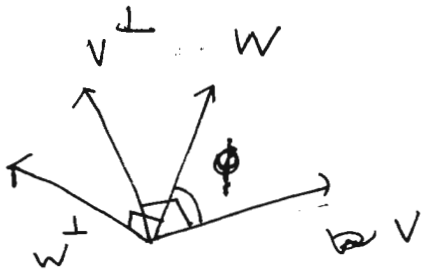
$$\left[\frac{Dw}{dt} \right] - \left[\frac{Dv}{dt} \right] = \frac{d\langle w, v \rangle}{dt} \quad (\langle \rangle \text{ is cancelled out})$$

hw: chapter 4:

4.2: 4, 6, 14, 17, 19

4.3: 1, 2, 3, 5, 7

4.4: 1, 2, 5, 12, 13, 14, 15, 17, 21, 23



Trivial Lemma

$$\cos \phi = w \cdot v$$

$$-\sin \phi \cdot \phi' = w' \cdot v + w \cdot v'$$

$$= -\sin \phi \cdot \lambda + \sin \phi \cdot \mu$$

so $\phi' = \lambda - \mu$

ie. $\left(\frac{Dw}{dt}\right) - \left(\frac{Dv}{dt}\right) = \frac{d\phi}{dt}$

now let $v = x_1 / |x_1| = e_1$ meaning of $\left[\frac{De_1}{dt}\right]$

extrinsic way:

pick low st $F=0$. ie. $v = e_1 = \frac{x_1}{\sqrt{E}}$, $e_2 = \frac{x_2}{\sqrt{G}}$

$$\star \frac{Dv}{dt} = \frac{D}{dt} \left(\frac{x_1}{\sqrt{E}} \right) = \frac{-1}{2\sqrt{E}} \left(\epsilon_1 \frac{du}{dt} + \epsilon_2 \frac{dv}{dt} \right) x_1$$

$$+ \frac{1}{\sqrt{E}} \left(x_{11}'' \frac{du}{dt} + x_{12}'' \frac{dv}{dt} \right)$$

$$\left[\frac{De_1}{dt} \right] = \frac{D}{dt} \left(\frac{x_1}{\sqrt{E}} \right) \cdot \frac{x_2}{\sqrt{G}} = \frac{1}{\sqrt{EG}} \left(x_{11} \cdot x_2 \frac{du}{dt} + x_{12} \cdot x_2 \frac{dv}{dt} \right)$$

$$\begin{matrix} (x_1 \cdot x_2)'_1 & - & x_1 \cdot x_{12} \\ \parallel & & \parallel \\ 0 & - & \frac{1}{2} \epsilon_2 \end{matrix} \quad \frac{1}{2} G_1$$

$$= \frac{1}{2} \frac{1}{\sqrt{EG}} \left(G_1 \frac{dv}{dt} - \epsilon_2 \frac{du}{dt} \right)$$

*: letter way (intrinsic)

$$\frac{1}{2\sqrt{E}} \left(\epsilon_1 \frac{du}{dt} + \epsilon_2 \frac{dv}{dt} \right) x_1 + \frac{1}{\sqrt{E}}$$

finally $\cdot \frac{x_2}{\sqrt{G}}$

$$= \frac{1}{\sqrt{EG}} \left(\Gamma_{11}^2 G u' + \Gamma_{21}^2 G v' \right)$$

$$\Gamma_{11}^2 = \frac{1}{2} g^{22} (\cancel{\partial_1 g_{12}} + \cancel{\partial_1 g_{12}} - \partial_2 g_{11})$$

$$= -\frac{1}{2} \epsilon_2 G^{-1}$$

$$\Gamma_{21}^2 = \frac{1}{2} g^{22} (\cancel{\partial_2 g_{11}} + \partial_1 g_{22} - \cancel{\partial_2 g_{21}})$$

$$= \frac{1}{2} G_1 G^{-1}$$

$$\frac{1}{G} = x_1 u' + x_2 v'$$

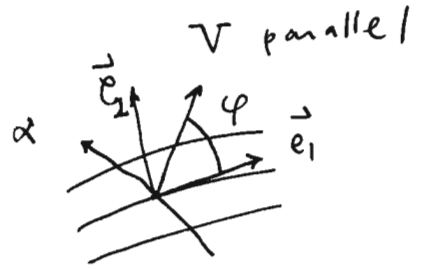
$$D_{\partial_t} x_1 = u' (\Gamma_{11}^1 x_1 + \Gamma_{11}^2 x_2)$$

$$+ v' (\Gamma_{21}^1 x_1 + \Gamma_{21}^2 x_2)$$

want only x_2 component.

$$\begin{pmatrix} E & 0 \\ 0 & G \end{pmatrix} \frac{1}{EG} \begin{pmatrix} +G & 0 \\ 0 & +E \end{pmatrix} = \begin{pmatrix} \frac{1}{E} & 0 \\ 0 & \frac{1}{G} \end{pmatrix}$$

$$\begin{aligned} \textcircled{1} \int_{\Omega} k dA &= \int_{\Omega} -\frac{1}{2} \frac{1}{\sqrt{EG}} \left(\left(\frac{E_2}{\sqrt{EG_2}} \right)_2 + \left(\frac{G_1}{\sqrt{EG}} \right)_1 \right) \sqrt{EG} du dv \\ &= \int_{\partial\Omega} \frac{1}{2} \left(\frac{E_2}{\sqrt{EG}} du + \frac{G_1}{\sqrt{EG}} dv \right) \\ &= - \int_{\partial\Omega} \left[\frac{d\vec{e}_1}{dt} \right] dt \\ &= \int_{\partial\Omega} \varphi' dt = \text{holonomy angle} \end{aligned}$$



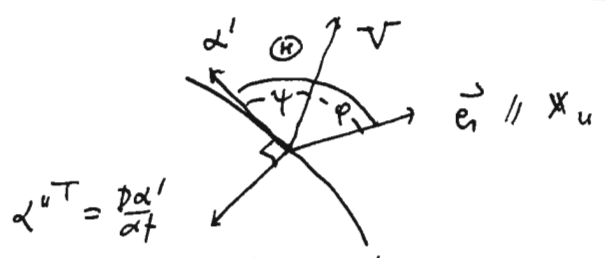
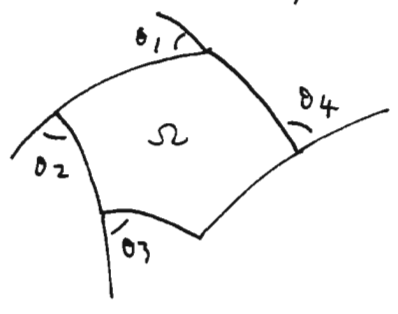
$$\textcircled{2} \varphi = \textcircled{H} - \psi$$

$\angle \text{angle}(\alpha', V)$
 $\angle \text{angle}(\alpha', \vec{e}_1)$

$$\begin{aligned} \cos \varphi &= V \cdot \vec{e}_1 \\ -\sin \varphi \cdot \varphi' &= V \cdot \lambda \vec{e}_2 \\ &= \sin \varphi \cdot \lambda \\ \Rightarrow \lambda &= -\varphi' \end{aligned}$$

$$\text{So } \int_{\partial\Omega} d\varphi = \int_{\partial\Omega} d\textcircled{H} - \int_{\partial\Omega} d\psi$$

$$\text{(Hopf)} \quad 2\pi - \sum_i \theta_i \quad \int_{\partial\Omega} kg ds$$



$$\begin{aligned} \alpha''^T &= \frac{d\alpha'}{dt} \\ \cos \psi &= \alpha' \cdot Z \\ -\sin \psi \cdot \psi' &= \alpha'' \cdot Z \\ &= kg \vec{n} \cdot Z \\ \cos(\psi + \frac{\pi}{2}) &= -\sin \psi \\ \Rightarrow \psi' &= kg \end{aligned}$$

ie. $\int_{\Omega} k dA + \int_{\partial\Omega} kg ds + \sum \theta_i = 2\pi$.
(local Gauss Bonnet)

here we pick t.s arc length for $\alpha = \partial\Omega$.

$$\textcircled{3} R \subset S \text{ oriented, } \partial R = \cup C_i$$

$$R = \bigcup_{j=1}^F T_j \text{ triangulation, } \theta_{ji}, i=1,2,3 \text{ outer angle}$$

$$\text{Sum } \Rightarrow \int_R k dA + \int_{\partial R} kg ds + \sum \theta_{jk} = 2F\pi$$

∂R closed

• 2 ways to interpret $3F$: $3\pi F - \sum \varphi_{jk}$ - interior angle

$$= \pi (2E_i + E_e) - (2\pi V_i + \pi V_{e,T} + \sum \theta_{jk})$$

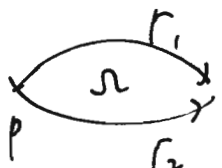
int' angle ext. angle from T ext. angle from C_i

$$\begin{aligned} &= 2\pi(E_i + E_e) - 2\pi(V_i + V_e) + \sum \theta_{jk} \\ &= 2\pi(E - V) + \sum \theta_{jk} \Rightarrow \text{global Gauss-Bonnet} \end{aligned}$$

Applications of Gauss Bonnet:

The 1st one is # ① in next page.


① cpt surface with $K \geq 0$ and not all $= 0$ is homeo to S^2 . (since $\chi = 2 - 2g > 0$)

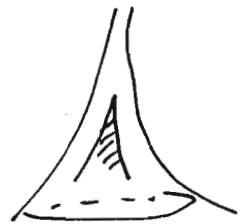
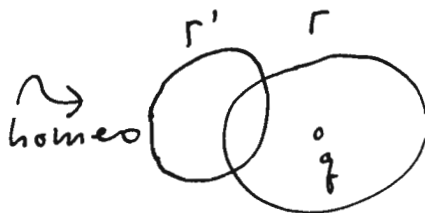
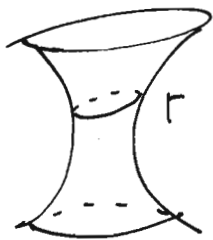
②  if $K \leq 0$ then no Γ_1, Γ_2 2 geodesics can bound a simple region Ω .

pf: $\int_{\Omega} K dA + \theta_1 + \theta_2 = 2\pi \Rightarrow \theta_1 = \theta_2 = \pi$ *

this also \Rightarrow no simple closed geodesic.
 $= \partial \Omega$. Ω simple.

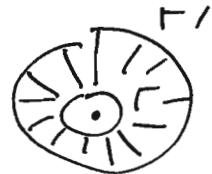
③ $S \sim$ homeo  cylinder with $K < 0$

$\Rightarrow \exists$ at most one simple closed geodesic  not this model



another Γ : if intersect Γ , must in 2 pts * to ②, it must be the case

Again $\int_{\Omega} K = 2\pi \chi(\Omega) = 0$ *



(No * if $K \equiv 0$) \rightarrow if S is closed in \mathbb{R}^3 (i.e. complete before, then will see $\exists!$ such Γ) minimize generator of $\pi_1(S) \cong \mathbb{Z}$.

④ S cpt, $K > 0$ Any 2 simple closed geodesics $\Gamma_1 \cap \Gamma_2 \neq \emptyset$.



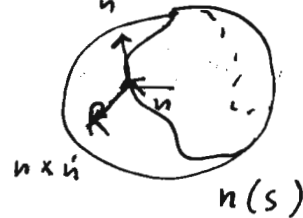
outside part Ω has $\chi(\Omega) = 0$

$\int_{\Omega} K dA = 2\pi \chi(\Omega) = 0$ *

this holds for any Ω orientable cpt surface with boundary

⑤ Jacobi's theorem:

$\alpha: I \rightarrow \mathbb{R}^3, K > 0$ if $n(s) \in S^2$ is a simple curve then n separates S^2 into 2 regions of equal area.



pf: \bar{S} arc length of α
 S arc length of n

$\text{Area}(R) = \int_R K dA = 2\pi - \int_{\partial R} kg d\bar{s}$

so this is equiv to show that $\int_{\partial R} kg d\bar{s} = 0$

$kg = \left[\frac{d\ddot{n}}{d\bar{s}} \right] = \ddot{n} \cdot (n \times \dot{n})$

in fact

$\ddot{n} = kg n \times \dot{n} + \text{normal part}$

$\dot{n} = n' \cdot \frac{ds}{d\bar{s}} = (-kT - \tau B) \frac{ds}{d\bar{s}}$

$\ddot{n} = (-kT - \tau B) \frac{d^2s}{d\bar{s}^2} + \underbrace{(-k'T - \tau'B)}_2 \left(\frac{ds}{d\bar{s}} \right) + (-k k \underline{n} - \tau \tau \underline{n}) \left(\frac{ds}{d\bar{s}} \right)^2$

by def of \bar{s} get $|\dot{n}| = 1$ hence $\frac{ds}{d\bar{s}} = \frac{k}{k^2 + \tau^2}$

$\Rightarrow kg = \ddot{n} \cdot (kB - \tau T) \frac{ds}{d\bar{s}}$

$= \cancel{(kT - kT)} \frac{ds}{d\bar{s}} \frac{ds}{d\bar{s}} + (k'T - kT') \cdot \left(\frac{ds}{d\bar{s}} \right)^3 + 0$

$= \frac{k'T - kT'}{k^2 + \tau^2} \cdot \frac{ds}{d\bar{s}} \quad (\text{notice } k = k' \cdot \frac{ds}{d\bar{s}})$

so $kg d\bar{s} = \frac{k'T - kT'}{k^2 + \tau^2} ds = -\tan^{-1} \left(\frac{\tau}{k} \right)' ds$

This should be the 1st one: thus $\int_{\partial R} kg d\bar{s} = -\tan^{-1} \frac{\tau}{k} \Big|_S = 0$.

⑥ Geodesic triangle



θ_i : outer angle

$\varphi_i = \pi - \theta_i$
inner angle

$\int_{\Omega} K dA + \sum_{i=1}^3 \theta_i = 2\pi$

ie. $-\pi + \sum_{i=1}^3 \varphi_i = \int_{\Omega} K dA$

- $\sum \varphi_i = \pi$ if $K = 0$
- $> \pi$ if $K > 0$
- $< \pi$ if $K < 0$

↑ the excess (angle) of Ω

= area of $N(\Omega)$ (in S^2)

should also explain this in terms of holonomy angle.

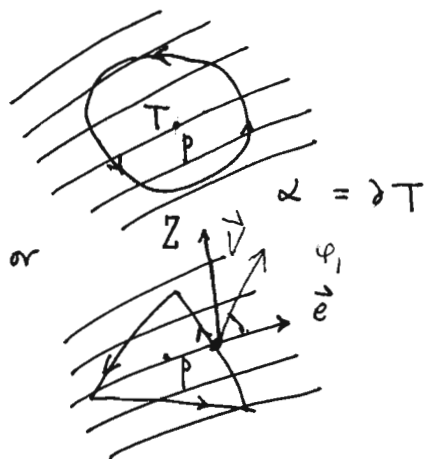
Hopf-Poincaré index thm

S cpt surface, \vec{v} sm tangent v.f. with isolated zeros
 ? orientable (or called sing. pts) \Rightarrow finite #

Defⁿ: (notation) index $I_p(\vec{v}) := \frac{1}{2\pi} \int_{\alpha} \varphi'_i \in \mathbb{Z}$

$$\varphi_i = \angle(\vec{e}, \vec{v})$$

indep of \vec{e} and α



Let $S = \cup T_i$, T_i contains at most one zero of \vec{v} .

$$\int_{T_i} K dA = \int_{\partial T_i} \varphi'_i$$

$\varphi = \angle(\vec{e}, Z)$ - parallel v.f. along ∂T_i , indep of Z

$$2\pi I_{p_i}(\vec{v}) = \int_{\partial T_i} \varphi'_i$$

$$\nRightarrow \int_S K dA - 2\pi \underbrace{\sum_i I_{p_i}(\vec{v})}_{\text{Index}(\vec{v})} = \sum_i \int_{\partial T_i} (\varphi - \varphi_i)' = 0$$

$$\varphi - \varphi_i = \angle(\vec{v}, Z)$$

$$\Rightarrow (\varphi - \varphi_i)' \text{ indep of } Z \quad *$$

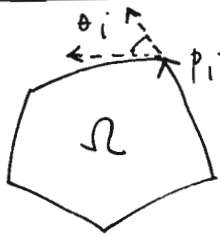
$$\text{So } 2\pi \text{Index}(\vec{v}) = \int_S K dA$$

$$\text{indep of } \vec{v} !! = 2\pi \chi(S)$$

$$\nRightarrow \boxed{\text{Ind}(\vec{v}) = \chi(S)}$$

Gauss Bonnet Theorem

Do Carmo's version:



$$\int_{\Omega} K dA$$

$$K := -\frac{1}{2\sqrt{EG}} \left[\left(\frac{E_2}{\sqrt{EG}} \right)_2 + \left(\frac{G_1}{\sqrt{EG}} \right)_1 \right]$$

$$= -\frac{1}{2} \int_{\Omega} \left[\left(\frac{G_1}{\sqrt{EG}} \right)_1 - \left(-\left(\frac{E_2}{\sqrt{EG}} \right)_2 \right) \right] du dv$$

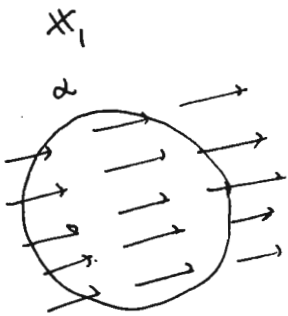
$$= -\frac{1}{2} \int_{\partial\Omega} -\frac{E_2}{\sqrt{EG}} du + \frac{G_1}{\sqrt{EG}} dv$$

$$= \int_{\partial\Omega} \left\{ \frac{1}{2} \left(\frac{E_2}{\sqrt{EG}} \right)_2 + \left(\frac{G_1}{\sqrt{EG}} \right)_1 \right\} ds$$

$$= \int_{\partial\Omega} \kappa_g ds - \left(\frac{DW}{dt} \right) dt$$

$$= \int_{\partial\Omega} - \left[\frac{D\vec{e}_1}{dt} \right] dt$$

geodesic curvature κ_g if $w = \alpha'$.



φ angle from x_1 to α' .
Hopf's theorem

$$\int_{\partial\Omega} d\varphi = 2\pi - \sum_i \theta_i$$

$$\Rightarrow \int_{\Omega} K dA + \int_{\partial\Omega} \kappa_g ds + \sum_{p_i} \theta_i = 2\pi$$

Local Gauss-Bonnet

$$D_{\alpha'} w \equiv \frac{Dw}{dt}$$

covariant diff.

Under the restriction that $|w|=1$.

then w^\perp is uniquely defined.

$$\text{so } \frac{Dw}{dt} = \lambda w^\perp \quad \text{i.e. } \lambda = \pm \left| \frac{Dw}{dt} \right| = \dot{w} \cdot w^\perp$$

call it $\left[\frac{Dw}{dt} \right]$ i.e. with sign

• when $w = \alpha'(s)$, s : arc length

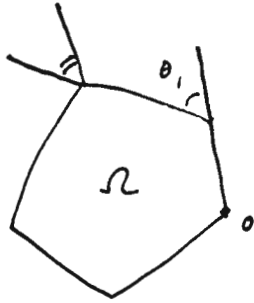
call $\kappa_g = \left[\frac{D\alpha'}{ds} \right]$ the geodesic curv. "algebraic value of $\frac{Dw}{dt}$ "

Fix any C^∞ field \vec{F} : let $\varphi(t) = \text{angle}(\vec{F}, w)$

Gauss-Bonnet Formula. (H. Hopf's version)

Thm: $\int_{\Omega} K dA + \int_{\partial\Omega} kg ds + \sum_i \theta_i = 2\pi$

Remark:

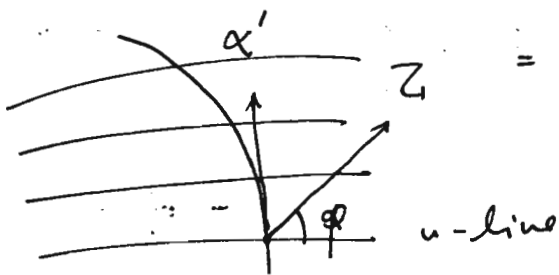


$2\pi - \sum \theta_i$ is exactly the holonomy angle.

" local Gauss-Bonnet \equiv holonomy def of K "

Pf: $ds^2 = dr^2 + G^2 d\theta^2$
 then $K = -\frac{G_{rr}}{G}$

$$\int_{\Omega} K dA = \int_{\Omega} -\frac{G_{rr}}{G} G dr d\theta = -\int_U G_{rr} dr d\theta = -\int_{\partial U} G r d\theta = \int_{\partial U} d\varphi$$



important formula for parallel translation

$\varphi = \text{angle of } Z \text{ to } du$
 $= \text{angle}(\alpha', du) - \text{angle}(\alpha', Z)$
 Z, α' both unit vector fields

$$\begin{aligned} \cos \varphi &= \alpha' \cdot Z \\ -\sin \varphi \cdot \varphi' &= \alpha'' \cdot Z + \alpha' \cdot Z' \\ &= kg \frac{\alpha' \cdot Z}{\cos(\varphi + \pi/2)} \Rightarrow \varphi' = kg \end{aligned}$$

geodesic curvature

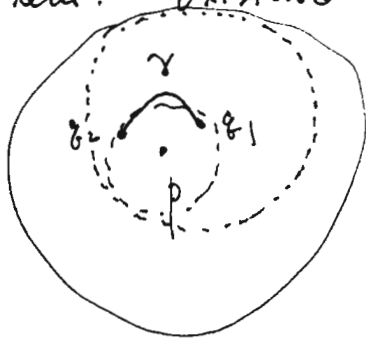
$$\int_{\Omega} K dA = \int_{\partial\Omega} d\varphi = \int_{\partial\Omega} d\theta - \int_{\partial\Omega} d\psi$$

holonomy | turning tangent via top. front = $2\pi - \sum \theta_i$ | geodesic curvature

Q.E.D.

Def¹⁷: distance on S : $d(p, \gamma) := \inf |\gamma|$

Theorem: Existence of convex neighborhood. γ joins p, q



ie. diff image of standard ball in $T_p S$ via exp $_p$

$B_r(p)$ good ball

$q \in \overline{B_{r/3}(p)}$ $B_{\delta_q}(q)$ good ball $\subset B$

$\frac{2}{3}r$

① why diffeo?

$$q' \in B_{\frac{2}{3}r}(q) \Rightarrow d(q', p) < \frac{r}{3} + \frac{2}{3}r \leq r.$$

then any $q_1, q_2 \in B_{r/3}(p)$, $d(q_1, q_2) < \frac{r}{3} + \frac{r}{3} = \frac{2r}{3}$

ie. $q_2 \in B_{2r/3}(q_1) \Rightarrow \exists!$ good γ joins q_1, q_2 . (shortest)

②

Claim: $B_{r/3}(p)$ is good. convex.

Ex.

diffeo -

for ①: $\overline{B_{r/3}(p)}$ cpt $\Rightarrow \exists \delta > 0$ st. $B_\delta(q)$ good ball $\forall q \in \overline{B_{r/3}(p)}$

so if $\delta < \frac{2}{3}r$, replace r by $\frac{3}{2}\delta$.

for ②: To show $\gamma \subset B_{r/3}(p)$, using normal coord (u^1, u^2)

$$f(t) := d^2(p, \gamma(t)) = \sum_i (u^i)^2$$

$$f'(t) = \sum_i 2u^i \dot{u}^i$$

$$f''(t) = 2 \sum_i (\dot{u}^i)^2 + 2 \sum_i u^i \ddot{u}^i$$

$$= 2 \left(\sum_i (\dot{u}^i)^2 - \sum_{i,j,k} \underbrace{u^i \Gamma_{jk}^i}_{\text{d}} \dot{u}^j \dot{u}^k \right)$$

Now pick r small enough such that $\sum_i |u^i \Gamma_{jk}^i| < \frac{1}{2} \forall j,k$

then $f''(t) \geq (\dot{u}^1 - \dot{u}^2)^2 \geq 0$

conti, = 0 at p

ie. f is convex, max occur at q_1 or q_2 done \square .

Corollary (EX). For S cpt, \exists triangulation by geodesic triangles.

hint: If $B_\delta(p)$ is convex $\forall p \in S$

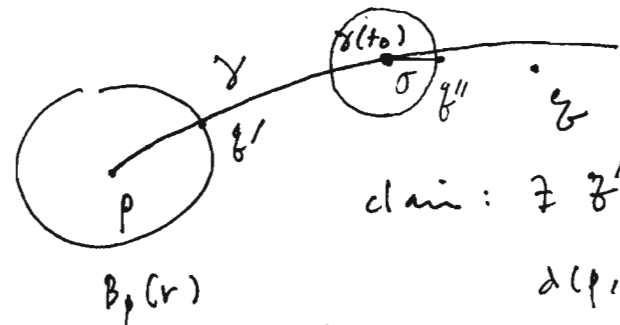
cover S by $B_{\delta/2}(p)$.



Lemma: $\exists p, \exp_p$ defined on $T_p M$

\Rightarrow any q is conn to p by a shortest geod.

pf:



$q \in B_p(r)$ done.

for $q \notin B_p(r)$,

claim: $\exists q' \in \partial B_p(r)$ st.

$$d(p, q) = r + d(q', q)$$

$B_p(r)$
normal cov ball
i.e. \exp_p is diffeo.

bf: Δ -ineq \Rightarrow " \leq " $\forall q'$

now let $d(q', q) = \inf_{q'' \in \partial B_p(r)} d(q'', q)$

let $\gamma: [0, \infty) \rightarrow M$

be the unique geod γ from p to q'

by def of d \nexists " \geq " $*$

$q'' \in \partial B_p(r)$
 \uparrow
 C^0 in q''

Continuity Method: $Z := \{t \mid d(p, \gamma(t)) + d(\gamma(t), q) = d(p, q)\}$

$Z \ni r (\neq \emptyset)$, closed, let $t_0 = \sup Z$, claim $\gamma(t_0) = q$.

otherwise, $\exists r_1 > 0$,

consider $\partial B_{\gamma(t_0)}(r_1) \ni q''$ st. $d(\gamma(t_0), q'') + d(q'', q) = d(\gamma(t_0), q)$

let σ the unique geod $\gamma(t_0) \rightarrow q''$.

$$\Rightarrow \underbrace{d(p, \gamma(t_0)) + d(\gamma(t_0), q)}_{d(p, q)} = \underbrace{d(p, \gamma(t_0)) + d(\gamma(t_0), q'') + d(q'', q)}_{d(p, q'') + d(q'', q)}$$

$$\Rightarrow d(p, \gamma(t_0)) + d(\gamma(t_0), q'') = d(p, q'') \quad (\leq \text{ by } *)$$

but then $\gamma|_{[0, t_0]} \cup \sigma|_{[0, r_1]}$ is a min. geod

$\Rightarrow \sigma \equiv \gamma$ and so $t_0 + r_1 \in Z$ $*$

Then (Hopf-Rinow) $\overset{1)}{\exists p, \exp_p} \Leftrightarrow \overset{2)}{\text{complete}} \Leftrightarrow \overset{3)}{\forall p, \exp_p}$

pf: $1) \Rightarrow 2)$ let $\{t_i\}$ Cauchy, $\exists \gamma_i(t_i) = q_i$, $t_i \rightarrow t_0$
arc lengths, Cauchy
Also $\gamma_i'(0) \rightarrow v$ in subsequence in $B_1(0)$ lift in $T_p M$.

consider $\gamma(0) = p, \gamma'(0) = v$

ODE thm $\Rightarrow \gamma_{i_n} = \gamma_{i_n}(t_{i_n}) \rightarrow \gamma(t_0)$, $\Rightarrow \gamma_i \rightarrow \gamma(t_0)$.

$2) \Rightarrow 3)$. If $\gamma(0) = p, \gamma'(0) = v$ only defined on $[0, t_0)$,

then $t_i \nearrow t_0 \Rightarrow \gamma(t_i)$ Cauchy $\Rightarrow \exists \gamma(t_0) = q$ extends $*$

Thm: $S \subset \mathbb{R}^3$ cpt $K = \text{const} \Rightarrow S \equiv \text{sphere}$.

let $k_1 \geq k_2$ C^∞ outside umbilical pts, C^0 on S .

lemma: If $K(p) > 0$, k_1 local max
 k_2 local min at p

then p is umbilical.

pf: If NOT, \exists curve (u, v) at p via line of curvature

hence $F = f = 0$, and $k_1 = \frac{f}{E} > k_2 = \frac{g}{G}$ (say)

since $e_2 = -dN(x_1) \cdot x_1 = k_1 \|x_1\|^2$

Recall (Mainardi-) Codazzi equations, in this case:

$$e_2 - \cancel{f_1} = e \frac{\Gamma_{12}^1}{11} + \cancel{f} (\Gamma_{12}^2 - \Gamma_{11}^1) - g \frac{\Gamma_{11}^2}{11} - \frac{1}{2} \frac{E_2}{G}$$

$$\frac{1}{2} g'' (\cancel{\partial_1 g_{21}} + \partial_2 g_{11} - \cancel{\partial_1 g_{21}}) = \frac{1}{2} \frac{G}{EG} E_2 = \frac{1}{2} \frac{E_2}{E}$$

i.e. $\begin{cases} e_2 = \frac{E_2}{2} \left(\frac{f}{E} + \frac{g}{G} \right), \text{ and by symmetry,} \\ g_1 = \frac{G_1}{2} \left(\frac{f}{E} + \frac{g}{G} \right). \end{cases}$

From $e = k_1 E \Rightarrow e_2 = (k_1)_2 E + k_1 E_2$

Also, $e_2 = \frac{E_2}{2} (k_1 + k_2)$

$\Rightarrow (k_1)_2 E = \frac{E_2}{2} (k_2 - k_1)$, symmetry

$(k_2)_1 G = \frac{G_1}{2} (k_1 - k_2)$.

Now, $K = -\frac{1}{2\sqrt{EG}} \left[\left(\frac{E_2}{\sqrt{EG}} \right)_2 + \left(\frac{G_1}{\sqrt{EG}} \right)_1 \right]$

$= -\frac{E_{22} + G_{11}}{2EG} + \frac{1}{4} \frac{(EG)_2}{(EG)^2} E_2 + \frac{1}{4} \frac{(EG)_1}{(EG)^2} G_1$

Key Point: $E_2 = \frac{2E}{k_2 - k_1} (k_1)_2 = * (k_1)_2$

$E_{22} = \frac{2E}{k_2 - k_1} (k_1)_{22} + * (k_1)_2$

* means some C^∞ function $\neq 0$ at p .

Hence, by symmetry

$$-2EG \cdot K = \frac{2E}{k_2 - k_1} (k_1)_{22} + \frac{2G}{k_1 - k_2} (k_2)_{11} + * (k_1)_2 + * (k_2)_1$$

$$\text{At } p: \quad \begin{array}{cccc} \wedge & & \vee & & \vee & & \parallel & & \parallel \\ & 0 & & 0 & & 0 & & 0 & & 0 \end{array}$$

get contradiction! \square

pf of Thm: We must have $K > 0$ (since \exists elliptic pt)

let p be the max of k_1 on S (p exists since S cpt)

then p is min of k_2 since $k_1 k_2 = K = \text{const}$.

$$\text{Lemma} \Rightarrow k_1(p) = k_2(p)$$

but then $\forall q \in S, k_1(q) \leq k_1(p) = k_2(p) \leq k_2(q)$

Hence $k_1(q) = k_2(q)$, i.e. All pts on S are umbilical.

$\Rightarrow S$ is a portion of a sphere, hence a sphere. \square

Thm': $S \subset \mathbb{R}^3$ cpt, $K > 0, H = \text{const} \Rightarrow S = \text{sphere}$.

↳ this means $S = \partial\Omega, \Omega$ convex
(called "ovaloid")

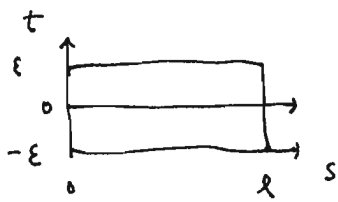
[Same Proof] .

Thm'' (Alexander-Hopf), $K > 0$ is NOT needed.

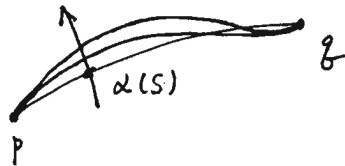
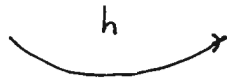
Defⁿ: $\alpha: [0, \ell] \rightarrow S$ by arc length s

variation $h: [0, \ell] \times (-\varepsilon, \varepsilon) \rightarrow S$ st $h_0 = \alpha$

s t $h_t(s) := h(s, t)$
 proper (end pt fixed) if $h_t(0) = \alpha(0), h_t(\ell) = \alpha(\ell)$



like a cov system on S , but h is only C^0 .
 h, dh may not be inj.



Fact: $V(s) := dh \begin{pmatrix} 0 \\ 1 \end{pmatrix}_{t=0} \equiv \frac{\partial h}{\partial t}(s, 0)$ is called the var. v.f.

$h'_t(s) \equiv T(s) := dh \begin{pmatrix} 1 \\ 0 \end{pmatrix} \equiv \frac{\partial h}{\partial s}(s, t)$ is the tangent v.f. of h_t

Any $V \in C^\infty$ on α , \exists variation h ,

if $V(0) = V(\ell) = 0$, h can be chosen to be proper

Pf: Consider $h(s, t) := \exp_{\alpha(s)}(tV(s))$. \square

Let $L(t) = \int_0^\ell \langle T, T \rangle^{1/2} ds = \text{length for } h_t$
 s is arc length only for $t=0$.

Thm: (1st Variation formula)

$$L'(0) = \langle V, T \rangle \Big|_0^\ell - \int_0^\ell \langle \nabla_T T, V \rangle ds \quad (\text{Here } T = \alpha')$$

Hence $L'(0) = 0 \forall$ proper variation $\iff \alpha$ is a geodesic.

$$\text{Pf: } L'(t) = \int_0^\ell \frac{d}{dt} \langle T, T \rangle^{1/2} ds = \int_0^\ell \frac{\langle \frac{\partial T}{\partial t}, T \rangle}{\sqrt{\langle T, T \rangle}} ds$$

Key point: $\frac{\partial T}{\partial t} = \nabla_V T = \nabla_T V$.

$$\Rightarrow \langle \frac{\partial V}{\partial s}, T \rangle = \frac{d}{ds} \langle V, T \rangle - \langle V, \frac{\partial T}{\partial s} \rangle$$

at $t=0, |T|=1$, get Thm. \square

time for cov system

$X = h, T = X_1, V = X_2$,

Also for extrinsic reason

$$\frac{\partial T}{\partial t} = \left(\frac{\partial}{\partial t} \frac{\partial h}{\partial s} \right)^T = \left(\frac{\partial}{\partial s} \frac{\partial h}{\partial t} \right)^T = \frac{\partial V}{\partial s}$$

2nd variation of geodesic.

$$L'(t) = \int_0^L \frac{\langle \partial_v T, T \rangle}{\langle T, T \rangle^{1/2}} ds$$

α : geodesic
 $v \perp T$ normal variation
 end-pt fixed

$$L''(t) = \int_0^L \left(\frac{\partial}{\partial t} \langle \partial_v T, T \rangle - \langle \partial_v T, T \rangle^2 \right) ds$$

$$+ \langle \partial_v \partial_T v, T \rangle + \langle \partial_v T, \partial_v T \rangle$$

$$= \int_0^L \left(|\partial_T v|^2 + \langle (\partial_v \partial_T - \partial_T \partial_v) v, T \rangle \right) ds$$

$$+ \langle \partial_T \partial_v v, T \rangle - \langle \partial_T v, T \rangle^2$$

$$+ \langle \partial_v v, \partial_T T \rangle - \langle v, \partial_T T \rangle^2$$

$$= \left. \langle \partial_v v, T \rangle \right|_0^L + \int_0^L \left(|\partial_T v|^2 - \langle [\partial_T, \partial_v] v, T \rangle \right) ds$$

Lemma: $\langle [\partial_T, \partial_v] v, T \rangle = K \cdot |v|^2$ (in general $|v \wedge T|^2$)
 the defect of commuting double covariant deri.

Pf: (How do we prove this in calculus on \mathbb{R}^2 ?)

A direct calculation by Gauss' equation. (uninteresting)

Actually it is infinitesimal version of G-B (i) via parallel tr.

Defⁿ: A surface is geod complete if every geod can be extended to ∞ . i.e. \exp_p is defined on whole $T_p S$, $\forall p \in S$.

Thm (Hopf-Rinow): geod complete \Leftrightarrow complete
 $v_p \Leftrightarrow \exists$ an p

And all \Rightarrow geod convex, any p, q are joined by shortest geod.

Thm (Bonnet): S complete, $k \geq k > 0 \Rightarrow d \leq \frac{\pi}{\sqrt{k}}$, S cpt.

Pf: If $\exists p, q \in S$, $d = d(p, q) > \frac{\pi}{\sqrt{k}}$

let α be the geod joins p, q , $e_0 \perp \alpha'(0)$, $|e_0| = 1$

$e(s)$ parallel along α

Consider $V(s) = \sin \frac{\pi s}{\lambda} e(s)$

$$0 \leq L''(0) \leq \frac{1}{2} \left(\frac{\pi^2}{\lambda^2} - k \right) < 0 \quad \leftarrow$$

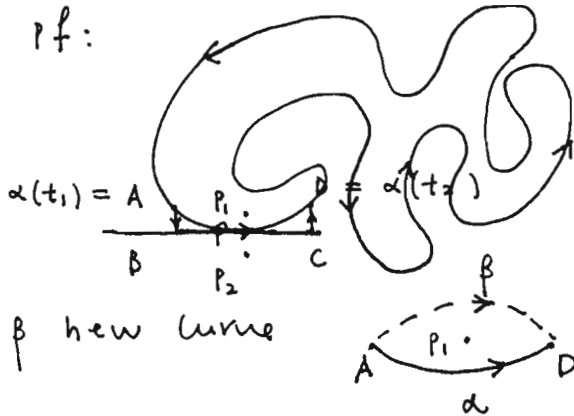


§ 5.7 Global Theory of Curves

(via degree of maps) $S \rightarrow S^1$ or $S \rightarrow S^2$

Thm 1 (C[∞] Jordan Curve Thm) $\deg \frac{\alpha - p}{|\alpha - p|}$ winding #
 piece-wise C² is OK.

Pf:



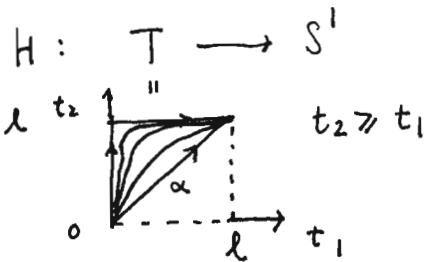
- $w(p)$ is constant in each conn comp of $\mathbb{R}^2 \setminus \alpha[0, l]$
- At least 2 comp $\Omega(p_1), \Omega(p_2)$: α is a graph near p
 $\Rightarrow w(p_2) - w(p_1) = \pm 1$ via β

- α is oriented $\Rightarrow \exists$ tubular nbd U , may set $p_1, p_2 \in U$
 - Ω any comp $\Rightarrow \partial_{top} \Omega \subset \alpha[0, l] \Rightarrow \Omega \cap \Omega(p_1) \neq \emptyset$ or $\Omega \cap \Omega(p_2) \neq \emptyset$.
- Thus \exists at most 2, hence exactly 2 conn. comp. \square

Thm 2 (Hopf's turning tangent)

α simple closed regular $C \mathbb{R}^2 \Rightarrow I_\alpha := \deg \alpha' = \pm 1$.
 rotation index

Pf: construct homotopy via secant lines:

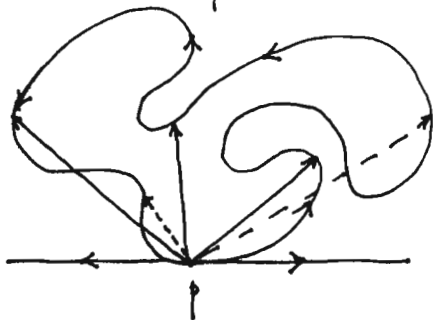


$$H(t_1, t_2) := \frac{\alpha(t_2) - \alpha(t_1)}{|\alpha(t_2) - \alpha(t_1)|}, \quad t_1 \neq t_2$$

$$H(t, t) := \frac{\alpha'(t)}{|\alpha'(t)|}$$

$$H(0, l) := -\frac{\alpha'(0)}{|\alpha'(0)|}$$

β new curve



H conti $\Rightarrow \deg \alpha = \deg \beta = \pm 1$. \square

Cor: A regular (closed) plane curve is convex \Leftrightarrow simple & k has sign.

\uparrow
 2 equiv def¹⁴, exterior vs interior

Pf (Ex + reading).

$\Leftrightarrow I = \pm 1$ and k has sign.

Space curves: (notice $k = |K(s)|$ for plane curves)

Thm 3 (Fenchel's thm) * $\int_{\alpha} k \geq 2\pi$, "=" $\Leftrightarrow \alpha$ plane convex.
 (simple) closed

p.f: Consider the tube of radius r : only need immersion
 $X(s, v) = \alpha(s) + r(\cos v \vec{n} + \sin v \vec{b})$, $[0, l] \times [0, 2\pi]$

Had seen: $X_s \times X_v = r(1 - rk \cos v) \cdot \vec{N}$

$$\vec{N} = -(\cos v \vec{n} + \sin v \vec{b})$$

$$N_s = k \cos v \vec{t} - \tau \sin v \vec{n} + b \cos v \vec{b}$$

$$N_v = \sin v \vec{n} - \cos v \vec{b}$$

$$\Rightarrow N_s \times N_v = k \cos v (\sin v \vec{b} + \cos v \vec{n}) = -k \cos v \vec{N}$$

Hence $K(s, v) = -\frac{k \cos v}{r(1 - rk \cos v)}$; $K \geq 0 \Leftrightarrow v \in [\frac{\pi}{2}, \frac{3\pi}{2}]$
 let $R = [0, l] \times [\frac{\pi}{2}, \frac{3\pi}{2}]$

$$\int_T K dA = 0 \quad (\text{useless})$$

But $\int_R K dA = -\int_0^l k(s) ds \int_{\pi/2}^{3\pi/2} \cos v dv = 2 \int_0^l k$

• Now any $N_0 \in S^2$ must be mapped by Gauss map over R
 (By the Far/Near parallel planes $\perp N_0$ argument)

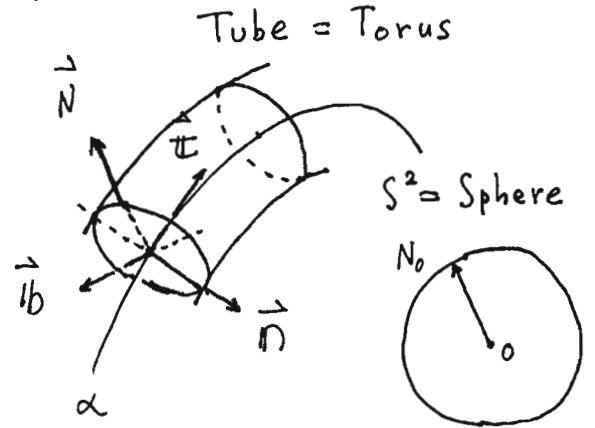
hence $2 \int_0^l k = \int_R K dA \geq 4\pi$.

• Moreover, for each s
 $T|_s \longrightarrow \Gamma_s \subset S^2$ great circle
 \cup
 $R|_s \longrightarrow \Gamma_s^+$

if α is a plane curve, $\partial \Gamma_s^+ = \{\pm P\}$ same $\forall s$ (north/south)

$\Gamma_{s_1}^+ \cap \Gamma_{s_2}^+$ only at $\pm P \Rightarrow \int_R K dA = 4\pi$.

* Rmk: For space curves $\int_{\alpha} k \in \mathbb{R}^+$ is in general not integer multiple of 2π .



Conversely, if $\int_{\alpha} k ds = 2\pi$ (so $\int_{\mathbb{R}} k dA = 4\pi$)

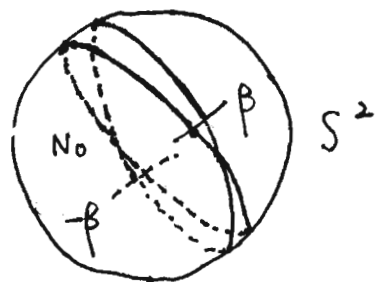
Claim:

$\Rightarrow \exists \mathbb{P} \in S^2$ st. $\partial \Gamma_s^+ = \{\pm \mathbb{P}\} \forall s$.

pf: if NOT, then $\partial \Gamma_s^+ = \{\pm \mathbb{P}(s)\}$ contains a curve β on S^2

Then $\exists s_1 \sim s_2$ with $\Gamma_{s_i} \cap \beta = \pm \mathbb{P}(s_i)$
in a nbd of $\partial \Gamma_{s_i}^+$. (ie. $\Gamma_{s_i} \neq \beta$)

Hence $\exists N_0 \in S^2$ being mapped by
2 points in $\mathbb{R}^+|_{s_1}$ & $\mathbb{R}^+|_{s_2}$ ($K > 0$)



The Gauss map is a local diffeomorphism there,

$\Rightarrow \int_{\mathbb{R}} k dA > 4\pi$ ~~*~~

Now, notice that $\frac{\overrightarrow{(\mathbb{P}(s))\mathbb{P}(s)}}{|\overrightarrow{(\mathbb{P}(s))\mathbb{P}(s)}} \parallel \vec{b}(s) \Rightarrow \vec{b} = \text{const} \Rightarrow \alpha$ plane curve.
 $\frac{\overrightarrow{(\mathbb{P}(s))\mathbb{P}(s)}}{|\overrightarrow{(\mathbb{P}(s))\mathbb{P}(s)}} \parallel \vec{b}(s)$ is a fixed direction why?

May assume $I_{\alpha} > 0$ (choose orientation) : (in fact = 1)

$$2\pi = \int_0^l |K| \geq \int_0^l K = 2\pi \cdot I_{\alpha} \geq 2\pi$$

Hence " \geq " holds, $K = |K| \geq 0 \Rightarrow$ convex. \square
"Simple" is not needed.

Thm 4 (Fary-Milnor) $\int_{\alpha} k ds < 4\pi \Rightarrow \alpha$ is unknotted.

pf: Choose a direction $v \in S^2$, $v \neq \vec{b}(s) \forall s \in [0, l]$

let $h_v(s) := \alpha(s) \cdot v$ the height fun wrt v .

$$h'_v(s) = \vec{t}(s) \cdot v$$

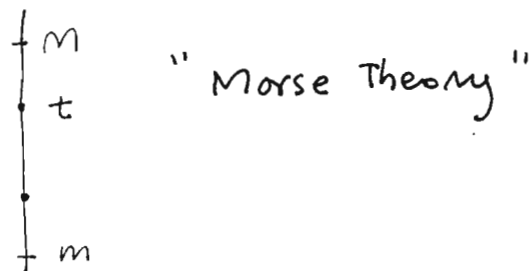
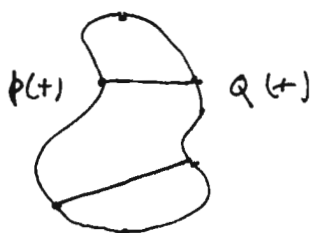
$$h''_v(s) = k \vec{n}(s) \cdot v$$

At any critical pt, we must have $\vec{n}(s) \cdot v \neq 0$

Since we may always assume $k > 0$,

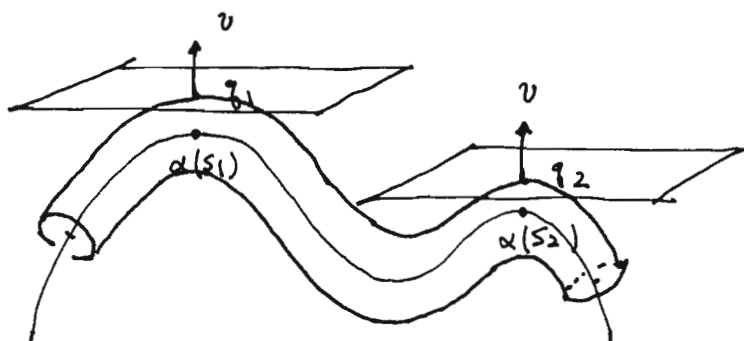
the critical pt must be local maxima or minima.

We claim that h_v has exactly 2 critical pts for some v and then the theorem is proved since the line segment joins $h_v^{-1}(t) = \{p(t), q(t)\}$



generate a surface $S \cong_{\text{homeo}} \text{disk}$, with $\partial S = \alpha$.

So, suppose h_v has ≥ 3 critical pts $\forall v \in \vec{b}([0, 2\pi])$
 WLOG, assume $h_v(s_1)$ max, $h_v(s_2)$ (local) max.



then $\exists \theta_1, \theta_2 \in \mathbb{R}$, $T_{\theta_1}, T_{\theta_2}$ has normal vector v
 $v \Rightarrow \vec{b}(s_1), \vec{b}(s_2) \Rightarrow$ in fact $K(\theta_1) > 0, K(\theta_2) > 0$.

But this works $\forall v \in \vec{b}([0, 2\pi])$, hence

$$\Rightarrow \int_{\mathbb{R}} K dA \geq 8\pi, \text{ i.e. } \int_{\alpha} k ds \geq 4\pi \quad \times \quad \square$$

Example of knots:

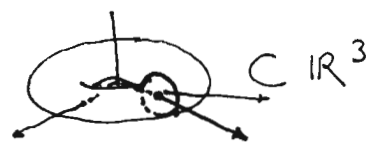
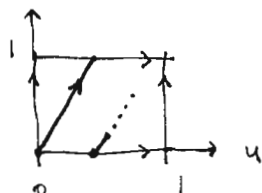
(1)



Trefoil knot
(simplest)
(2,3)

(2) Torus knot $(p, q) = 1$

v



$$t \in [0, 2\pi] \mapsto (u, v) \mapsto \begin{aligned} x &= (R + r \cos v) \cos u \\ y &= (R + r \cos v) \sin u \\ z &= r \sin v \end{aligned}$$

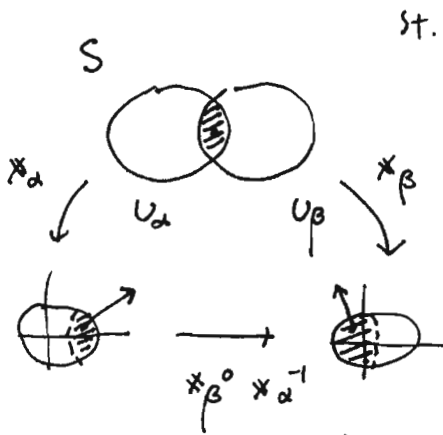
Can $\int k = 4\pi$ (why?)

Ex. $\int k = ?$

minimize it over all (R, r) .

Ans: Unless it approaches to the plane!

- An abstract surface (or 2-dim'l manifold) is a top space S covered by coord charts $(U_\alpha, \kappa_\alpha)$



st. $\phi_{\beta\alpha} := \kappa_\beta \circ \kappa_\alpha^{-1}$ is $C^k \forall \alpha, \beta, k \geq 0$.

equiv. $S = (\coprod_\alpha U_\alpha) / \sim$

- For $k \geq 1$, the tangent space is defined as $T_p U_\alpha$, for $p \in U_\alpha$, under

identification $d\phi_{\beta\alpha} : T_p U_\alpha \xrightarrow{\sim} T_p U_\beta$
 denote a basis by $\kappa_i \equiv \partial_i$ in $T_p U_\alpha$ (why use ∂_i ?)

- A Riemannian metric is a collection of 1st fund. form

$\langle \cdot, \cdot \rangle_\alpha$ on TU_α st. $\langle v, w \rangle_\alpha = \langle d\phi_{\beta\alpha}(v), d\phi_{\beta\alpha}(w) \rangle_\beta \forall \alpha, \beta$.

Denote as (S, g) . g always exists by P.O.U.

denote $g_{ij} = \langle \kappa_i, \kappa_j \rangle$, $\alpha' = \sum \kappa_i du^i \neq ds^2 = \langle \alpha', \alpha' \rangle$

- The covariant derivative (Levi-Civita connection) = $\sum g_{ij} du^i du^j$

∇_v : Vect fields \rightarrow Vect fields at $v \in T_p S$

is uniquely determined by, let $v = \alpha'$, $\nabla_v \equiv \frac{D}{dt}$:

1) linear op: $\nabla_v (a w_1 + b w_2) = a \nabla_v w_1 + b \nabla_v w_2, a, b \in \mathbb{R}$

$\nabla_{a v_1 + b v_2} w = a \nabla_{v_1} w + b \nabla_{v_2} w$

2) Leibnitz rule: $\nabla_v (f w) = (v f) w + f \nabla_v w$

L.C. condition:

3) Metrical: $\frac{d}{dt} \langle w_1, w_2 \rangle = \langle \nabla_v w_1, w_2 \rangle + \langle w_1, \nabla_v w_2 \rangle = \frac{d \langle \alpha', \alpha' \rangle}{dt}$ (directional derivative)

4) Torsion free: For $\nabla_{\kappa_i} \kappa_j =: \sum \Gamma_{ij}^k \kappa_k$
 $\nabla_{\kappa_i} \kappa_j = \nabla_{\kappa_j} \kappa_i$, i.e. $\Gamma_{ij}^k = \Gamma_{ji}^k$

Then $\Gamma_{ij}^k = \frac{1}{2} g^{kl} (\partial_i g_{jl} + \partial_j g_{il} - \partial_l g_{ij})$.

- K is determined by the Gauss eq'ns (any one of it)

\Rightarrow parallel translation

Gauss-Bonnet Thm

exp map, geod. convex nbd

variation of geodesics (Bonnet's thm etc)

Remark: For 2-dim'l reason, passing to the univ. cover

$$(\tilde{S}, \tilde{g}, \pi) \rightarrow (S, g) \quad \pi_1(\tilde{S}) = \{1\} \Rightarrow \begin{cases} \tilde{S} \cong S^2 & \text{cpt} \\ \tilde{S} \cong \mathbb{R}^2 & \text{non-cpt} \\ & \text{homeo.} \end{cases}$$

thus for $S \neq S^2, \mathbb{R}P^2$, may assume there is only one chart (U, κ) .

• Geometry of Space forms (S with $K = \text{const}$, complete) §5.5 + §5.6 + §5.8 (Cartan-Hadamard thm) Uniformization Thm:

Elliptic	$K = 1$	\Leftrightarrow	S isometric to S^2	$(S^2, \mathbb{R}P^2)$	$z=1$
Euclidean	$K = 0$	\Leftrightarrow	"	\mathbb{R}^2	$(\mathbb{C}, \mathbb{D}, \mathbb{M})$ cover
Hyperbolic	$K = -1$	\Leftrightarrow	\mathbb{H}	(\mathbb{H}/Γ)	discrete subgroup of isometries

Remark: Any cpt S can admit g with $K = \text{const}$.

Need PDE to prove (Kobayashi-Werner, or Riem mapping) by changing $g \mapsto fg$ ($f > 0, C^\infty$) i.e. conformal.

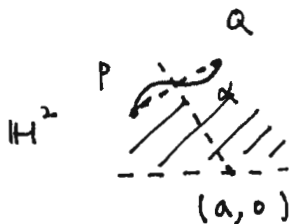
• Poincaré model of \mathbb{H} (4 models)



$$ds^2 = e^{-2y} dx^2 + dy^2 \quad \text{let } u=x, v=e^y$$

$$= \frac{du^2 + dv^2}{v^2} \quad \Rightarrow dv = v dy$$

& $v \in (0, \infty)$.



$$K = -\frac{1}{2\sqrt{EG}} \left[\left(\frac{E_2}{\sqrt{EG}} \right)_2 + \left(\frac{G_1}{\sqrt{EG}} \right)_1 \right] \equiv -1 \quad (\text{easy})$$



geodesic: $u = a + r \cos \theta$
 if $\vec{PA} \neq \vec{e}_2$: $v = r \sin \theta$

$$|\alpha| = \int_{\theta_1}^{\theta_2} \frac{\sqrt{dx^2 + r^2 d\theta^2}}{r \sin \theta} \geq \int_{\theta_1}^{\theta_2} \frac{d\theta}{\sin \theta}$$

"=" $\Leftrightarrow r = \text{const}$. i.e. α is a circle

Also, then $|\alpha| = -\ln(\csc \theta + \cot \theta) \Big|_{\theta_1}^{\theta_2} \rightarrow \infty$ as $\theta_1 \rightarrow \pi$
 i.e. \mathbb{H} is complete.

if $\vec{PA} \parallel \vec{e}_2$ get $d\theta = 0$.
 i.e. $\theta = \pi/2$

$$\mathbb{H} \rightarrow D, \quad z = \frac{w-i}{w+i}$$

$$(i, i, 0) \mapsto (1, 0, -1)$$

Ex. Determine geodesics in D & $\int_{\alpha} K dA$

CLAIM: indeed $ds^2 = \frac{4|dz|^2}{(1-|z|^2)^2}$ and $\text{isom}(D) = \text{Aut}(D) = \left\{ e^{i\theta} \frac{z-d}{1-\bar{d}z} \right\}$
 by Schwarz lemma.

$$z = \frac{w-i}{w+i} \Leftrightarrow$$

$$zw + zi = w - i$$

$$w(1-z) = i(1+z) \Leftrightarrow w = i \frac{1+z}{1-z}$$

$$y = \text{Im } w = \text{Re } \frac{1+z}{1-z} = \frac{1}{2} \left(\frac{1+z}{1-z} + \frac{1+\bar{z}}{1-\bar{z}} \right) = \frac{1}{2} \frac{2(1-|z|^2)}{|1-z|^2}$$

$$\begin{cases} w = u + iv \\ z = x + iy \end{cases} \quad z = \frac{aw+b}{cw+d} \Rightarrow dz = \frac{a(cw+d) - (aw+b)c}{(cw+d)^2} dw = \frac{ad-bc}{(cw+d)^2} dw$$

$$\Rightarrow |dz| = |ad-bc| \cdot |dw| \cdot \frac{1}{|cw+d|^2}$$

$$= 2 \frac{|dw|}{|w+i|^2}$$

$$= \frac{1}{2} |1-z|^2 |dw|$$

$$w+i = i \left(\frac{1+z}{1-z} + 1 \right) = \frac{2i}{1-z}$$

$$\text{i.e. } ds^2 = \frac{4|dz|^2}{(1-|z|^2)^2} \text{ as expected.}$$

- $\text{Aut}(D) = \left\{ e^{i\theta} \frac{z-a}{1-\bar{a}z} \mid a \in D \right\}$ by Schwartz lemma in complex analysis.
- $\text{Isom}(D) \supset \text{Aut}(D)$ by direct check.
- \subset since Isometric $\subset \text{Conf}^+(D) = \text{Aut}(D)$.

Hyperbolic projection:

$$t^2 - (x^2 + y^2) = 1$$

$$\begin{matrix} \text{!!} \\ p^2 \end{matrix} \quad t^2 = 1 + p^2 \Rightarrow t dt = p dp$$

$$ds^2 = dt^2 - dp^2 = \frac{p^2}{1+p^2} dp^2 - dp^2 = \frac{dp^2}{1+p^2}$$

$$\text{Also, } r = \frac{p}{t+1}$$

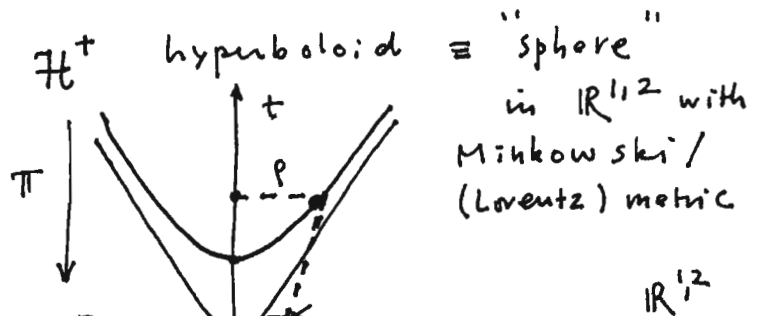
$$\Rightarrow p = rt + r$$

$$\Rightarrow (p-r)^2 = r^2 t^2 = r^2 (1+p^2)$$

$$\text{i.e. } p - 2r = pr^2$$

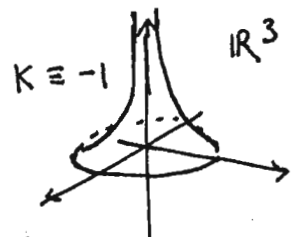
$$\Rightarrow \boxed{p = \frac{2r}{1-r^2}}$$

$$\text{So } ds^2 = \frac{\left[\frac{2(1-r^2) + 4r^2}{(1-r^2)^2} \right]^2 dr^2}{1 + \frac{4r^2}{(1-r^2)^2}} = \frac{4 dr^2}{(1-r^2)^2} \equiv \frac{4(dx^2 + dy^2)}{(1-x^2-y^2)^2}$$



Notice that the pseudosphere in \mathbb{R}^3 is NOT complete.

yet it is in $\mathbb{R}^{1,2}$.



Ex. $E \ni 0$ a plane $\Rightarrow C = E \cap \mathbb{H}^+$ geodesic $\Leftrightarrow \pi(C)$ a circle $\perp S^1 = \partial D$.

Ref: Dubrovin (Fomenko / Novikov : Modern Geometry I. Ch 2. (8/10)

Thm (Hilbert) S complete, $K = -1$

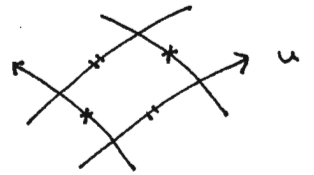
\Rightarrow ~~isometric~~ isometric immersion $S \rightarrow \mathbb{R}^3$.

Step 1: $K = -1 \Rightarrow \exists$ T-net
 pf: If $S \rightarrow \mathbb{R}^3$, asymptotic curves form a

~~Isobolous~~ T-net: i.e. asymp cov (u, v)

$$e(u')^2 + 2f u'v' + g(v')^2 = 0 \Rightarrow e = 0 = g$$

$X(u, v)$ should have $E_2 = 0 = G_1$ i.e. $X_{12} \perp N$



$$N_1 \times N_2 = K X_1 \times X_2 = K D N, \quad D := \sqrt{EG - F^2}$$

$$-1 = K = \frac{eg - f^2}{EG - F^2} = -\left(\frac{f}{D}\right)^2 \Rightarrow \frac{f}{D} = \pm 1$$

$$= |X_1 \times X_2|$$

opposite sides have equal length

$$N \times N_1 = \frac{1}{D} (X_1 \times X_2) \cdot N_1 = \frac{1}{D} [(X_1 \cdot N_1) X_2 - (X_2 \cdot N_1) X_1] = \frac{f}{D} X_1$$

-f = 0

Similarly $N \times N_2 = -\frac{f}{D} X_2$

$$(N \times N_1)_2 - (N \times N_2)_1 = \pm 2 X_{12} \Rightarrow X_{12} = \pm N$$

$$= 2 N_1 \times N_2 = 2 N$$

Step 2: for any T-net
curves in T-net

$$\left(\begin{array}{l} \text{so } E_2 = (X_1 \cdot X_1)_2 = 2 X_{12} \cdot X_1 = 0 \\ G_1 = 0 \text{ too.} \end{array} \right)$$

$$E_2 = 0 \Rightarrow \text{set } \tilde{u} = \int E du$$

may assume (u, v) st

$$G_1 = 0 \Rightarrow \text{set } \tilde{v} = \int G dv$$

$$ds^2 = du^2 + 2 \underbrace{\cos w}_{X_1 \cdot X_2} du dv + dv^2$$

§4.3 Ex.5 $\Rightarrow K = -\frac{w_{uv}}{\sin w}$

i.e. $w_{12} + K \sin w = 0$.

Corollary: In any rectangular region R_0 , $\left| \int_{R_0} K dA \right| < 2\pi$

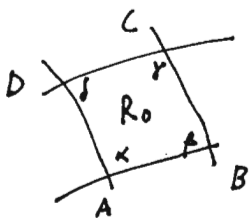
pf: $\int_{R_0} K dA = \int_{R_0} K |X_1 \times X_2| du dv = \int_{R_0} K \sin w du dv$

$$= - \int_{R_0} w_{uv} du dv$$

$$= - \int_B^C w_v dv + \int_A^B w_v dv$$

$$= -(\gamma - (\pi - \beta)) + (\pi - \delta) - \alpha = 2\pi - (\alpha + \beta + \gamma + \delta)$$

$$\in (-2\pi, 2\pi). \quad \square$$



Step 3: Global T-net

By taking the universal covering $\hat{S} \rightarrow S$

(in fact, $\exp_p: T_p S \cong \mathbb{R}^2 \rightarrow S$ is a covering map)

may assume that $S = \mathbb{H}^2$ (or \mathbb{D}), which has no vol.

isometric

by Minding's thm, since $K < 0 \neq \exp_p$ is defined on the whole $T_p S$

Now the asymptotic coord system is also globally defined on S . (Again using completeness of S)

$$\text{But then } S = \bigcup_{n=1}^{\infty} X([-n, n] \times [-n, n]) =: \bigcup_{n=1}^{\infty} Q_n$$

$$|Q_n| = - \int_{Q_n} K dA < 2\pi \text{ in step 2 } (K \equiv -1)$$

$$Q_n \subset Q_{n+1} \subset \dots \Rightarrow |S| < \infty \quad \times$$

Some Historical Remarks on isometric embeddings in \mathbb{R}^n :

Global

C^∞ topological: Whitney $M^n \hookrightarrow \mathbb{R}^{2n}$ so $S \hookrightarrow \mathbb{R}^4$

isometric: Nash (1956) Günther (1989)

$$M^n \hookrightarrow \mathbb{R}^N \quad N = \frac{3}{2}n(n+1)(n+9); \quad \frac{1}{2}n(n+3)+5$$

in any ϵ -ball

$$\text{so } S \hookrightarrow \mathbb{R}^{10}$$

For $S = \mathbb{H}^2$: Blaušá (1955), \exists explicit $\mathbb{H}^2 \hookrightarrow \mathbb{R}^6$
 C^∞ isometric

Local

Conjecture (Yau): $S \hookrightarrow \mathbb{R}^4$ isometric C^∞
for C^1 OK by Nash-Kuiper (1956).

~~Any local metric~~ $S \hookrightarrow \mathbb{R}^4$, isod. into $x^2+y^2 = z^2+w^2$

~~For \mathbb{R}^3~~ , OK near $p \in S$ if $K(p) \neq 0$

or if $K(p) = 0$ but $\nabla K(p) \neq 0$ (or finite order)
(C.S. Lin 1986).

for \mathbb{H}^2 certainly OK locally (Minding) ; globally in $\mathbb{R}^{1,2}$.