PROBLEMS FOR CALCULUS 2009 - FINAL EXAM

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Problem 0.1. Using Green's Theorem to derive the change of variable formula:

$$\iint_{\Omega} f \, dx \, dy = \iint_{D} f \circ T \left| \frac{\partial(x, y)}{\partial(u, v)} \right| \, du \, dv,$$

where the map $T : (u, v) \mapsto (x(u, v), y(u, v))$ has continuous partial derivatives and maps D one to one and onto Ω .

Solution. Let $Q(x,y) = \int_{-\infty}^{x} f(t,y) dt$. Then $Q_x = f$ and by Green's Theorem $\iint_{\Omega} f \, dx dy = \iint_{\Omega} Q_x \, dx dy = \int_{\partial \Omega} Q \, dy.$

By the change of variable formula for one variable (line) integrals, this equals

$$\pm \int_{\partial D} Q(y_u \, du + y_v \, dv) = \pm \int_{\partial D} (Qy_u) \, du + (Qy_v) \, dv$$

Here the sign depends on whether *T* preserves the orientation. By Green's Theorem again, now on the (u, v) plane, we get

$$\pm \iint_D \left((Qy_v)_u - (Qy_u)_v \right) du dv.$$

The integrand equals $(Q_x x_u + Q_y y_u)y_v + Qy_{vu} - (Q_x x_v + Q_y y_v)y_u - Qy_{uv}$, which simplifies to $Q_x(x_u y_v - x_v y_u) = f(x_u y_v - x_v y_u)$ by the definition of Q.

Finally, by the definition of orientation of T we have

$$\pm (x_u y_v - x_v y_u) = \left| \frac{\partial(x, y)}{\partial(u, v)} \right|$$

The proof is complete.

Problem 0.2. The Maxwell equations without charges and currents state that

$$\operatorname{div} ec{E} = 0 = \operatorname{div} ec{H}, \quad \operatorname{curl} ec{E} = -rac{1}{c} rac{\partial ec{H}}{\partial t}, \quad \operatorname{curl} ec{H} = rac{1}{c} rac{\partial ec{E}}{\partial t}.$$

Using them to prove the wave equation for the electronic field \vec{E} :

$$\nabla^2 \vec{E} = \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2}.$$

Solution. From

$$\operatorname{curl}\operatorname{curl}\vec{E} = \nabla\operatorname{div}\vec{E} - \nabla^2\vec{E} = -\nabla^2\vec{E}$$

(since div $\vec{E} = 0$), we get

$$\nabla^2 \vec{E} = -\operatorname{curl}\left(-\frac{1}{c}\frac{\partial \vec{H}}{\partial t}\right) = \frac{1}{c}\frac{\partial\operatorname{curl}\vec{H}}{\partial t} = \frac{1}{c^2}\frac{\partial^2 \vec{E}}{\partial t^2},$$

where in the middle equality we commute the derivatives in *t* and in *x*, *y*, *z*. \Box

Problem 0.3. Let \vec{F} be a conservative vector field on \mathbb{R}^3 with potential P(x, y, z). State and prove the "Law of Conservation of Energy".

Solution. For a moving particle $\vec{r}(t)$ of mass *m*, we need to show that

$$E(t) := \frac{1}{2}mv^2(t) + P(\vec{r}(t))$$

is a constant value in *t*.

Method 1: By differentiation in *t* and using $\vec{F} = -\nabla P$, we get

$$E'(t) = \frac{1}{2}m(\vec{v}\cdot\vec{v})' + \nabla P\cdot\vec{r}'(t)$$
$$= m\vec{a}\cdot\vec{v} - \vec{F}\cdot\vec{v} = \vec{F}\cdot\vec{v} - \vec{F}\cdot\vec{v} = 0$$

Thus E(t) is independent of t.

Method 2: By computing the work $W = \int_C \vec{F} \cdot d\vec{r}$ done along *C* in two ways:

$$W(a,b) = \int_{a}^{b} \vec{m}\vec{a} \cdot \vec{v} \, dt = \int_{a}^{b} \frac{1}{2}m(\vec{v} \cdot \vec{v})' \, dt = \frac{1}{2}mv^{2}(b) - \frac{1}{2}mv^{2}(a).$$

On the other hand, since $\vec{F} = -\nabla P$,

$$W(a,b) = -\int_{a}^{b} \nabla P \cdot \frac{d\vec{r}}{dt} \, dt = -\int_{a}^{b} P(\vec{r}(t))' \, dt = -P(\vec{r}(b)) + P(\vec{r}(a)).$$

The conservation follows by substraction of the above two expressions.

Problem 0.4. Solve $y'' - 2y' + y = e^x / (1 + x^2)$.

Solution. The characteristic equation for y'' - 2y' + y' = 0 is $(r - 1)^2 = 0$, hence the solutions for the homogenuous equation is given by

$$c_1e^x + c_2xe^x$$
.

By variation of parameters we consider a particulat solution as

$$y_p = ue^x + vxe^x.$$

Then $y'_p = u'e^x + ue^x + v'xe^x + ve^x + vxe^x$.

We "need to" impose $u'e^x + v'xe^x = 0$, that is u' + v'x = 0. Then

$$y'_p = ue^x + ve^x + vxe^x = y_p + ve^x$$

$$y_p'' - 2y_p' + y_p = (y_p + v'e^x + 2ve^x) - 2(y_p + ve^x) + y_p = v'e^x.$$

The equation to be solved becomes $v' = 1/(1 + x^2)$ and so $v = \tan^{-1} x$. Then $u' = -xv' = -x/(1 + x^2)$ and so $u = -\frac{1}{2}\ln(1 + x^2)$. So the general solutions are

$$y(x) = e^{x} \left(c_1 + c_2 x - \frac{1}{2} \ln(1 + x^2) + \tan^{-1} x \right).$$

Problem 0.5. Find all power series solutions to the (special case of) Bessel equation

$$x^2y'' + xy' + x^2y = 0$$

What can be concluded for the initial data y(0) and y'(0)?

Solution. Let $y(x) = \sum_{n \ge 0} a_n x^n$. The equation becomes

$$\sum_{n=2}^{\infty} a_n n(n-1)x^n + \sum_{n=1}^{\infty} a_n nx^n + \sum_{n=0}^{\infty} a_n x^{n+2} = 0.$$

There are no x^0 terms. For x^1 terms we get $a_1x = 0$ hence $a_1 = y'(0) = 0$. By comparing coefficients of x^k with $k \ge 2$ we get

$$a_k k(k-1) + a_k k + a_{k-2} = 0.$$

So

$$a_k = -\frac{a_{k-2}}{k^2}.$$

By iteration, if *k* is odd then we must have $a_k = 0$. For k = 2m we get

$$a_{2m} = (-1)^m \frac{1}{(2m)^2 \times (2(m-1))^2 \cdots \times 2^2} a_0 = \frac{(-1)^m}{2^{2m} (m!)^2} a_0.$$

So all the power series solutions are

$$y(x) = a_0 \sum_{m=0}^{\infty} \frac{(-1)^m}{2^{2m} (m!)^2} x^{2m}.$$

(The power series clearly onverges for all $x \in \mathbb{R}$.)

The initial condition $y(0) = a_0$ is arbitrary, but $y'(0) = a_1$ must be zero.